











Bulge+Disk Morphology in Rest-frame UV and Optical: Size–Mass Relations Reveal Distinct Growth Paths for Star-forming and Quiescent Galaxies

ANGELO GEORGE (喬安哲) ^{1,2} IVANA DAMJANOV ^{1,*} MARCIN SAWICKI ¹ STÉPHANE ARNOUITS ³ LINGJIAN CHEN,¹
GUILLAUME DESPREZ ¹ STEPHEN GWYN ⁴ THIBAUD MOUTARD ⁵ MICHELE PIZZARDO ^{1,6} ANNA SAJINA ⁷ AND
DEVIN J. WILLIAMS ¹

¹*Institute for Computational Astrophysics and Department of Astronomy & Physics, Saint Mary’s University, 923 Robie Street, Halifax, Nova Scotia, B3H 3C3, Canada*

²*Institute for Astronomy and Astrophysics, Academia Sinica, No.1, Sec. 4, Roosevelt Rd, Taipei 106319, Taiwan*

³*Aix Marseille Université, CNRS, CNES, LAM, 38 rue Frédéric Joliot-Curie, Marseille 13388 cedex 13, France*

⁴*NRC Herzberg Astronomy and Astrophysics, 5071 West Saanich Road, Victoria, British Columbia, V9E 2E7, Canada*

⁵*European Space Agency (ESA), European Space Astronomy Centre (ESAC), Camino Bajo del Castillo s/n, 28692 Villanueva de la Cañada, Madrid, Spain*

⁶*Institute for Astronomy and Astrophysics, Academia Sinica, Taipei 106319, Taiwan*

⁷*Department of Physics and Astronomy, Tufts University, Medford, MA 02155, USA*

ABSTRACT

We characterize the structural evolution of a mass-complete sample of 44,000 star-forming galaxies (SFGs) and 34,000 quiescent galaxies (QGs) at $0.1 < z < 0.85$ using the CFHT CLAUDS and Subaru HSC-SSP datasets. By performing bulge+disk decompositions in both rest-frame UV (3000 Å) and optical (5000 Å), we provide the first statistical dual-wavelength analysis of the size–mass relation for these individual components. At a fiducial mass of $M_* = 5 \times 10^{10} M_\odot$, disk sizes are consistently larger than their corresponding bulges, with this disparity being more pronounced in the UV. We find that both components are larger in SFGs than in QGs. Notably, disks are more extended in the UV than in the optical, while bulge sizes remain comparable across both wavelengths. The pace of size evolution (measured as $R_e \propto (1+z)^\beta$) is wavelength-dependent: QG disks grow significantly faster in the optical ($\beta = -0.99$) than in the UV ($\beta = -0.49$), whereas bulges exhibit the opposite trend, with faster growth in the UV ($\beta = -1.51$) than in the optical ($\beta = -1.02$). SFG components show similar but less pronounced trends, with disks evolving at $\beta = -0.53$ (UV) and -0.64 (optical), whereas bulges evolve at $\beta = -1.09$ (UV) and -0.70 (optical). Our results suggest that bulge growth primarily drives the size evolution of SFGs, while the structural assembly of QGs is governed by a combination of steady influx of recently quenched ‘newcomers’ and minor mergers/accretion.

Keywords: Galaxies (573) — Galaxy evolution (594) — Ultraviolet astronomy (1736) — Optical astronomy (1776)

1. INTRODUCTION

Galaxies assemble their stellar mass through a combination of processes — including star formation, mergers, gas accretion, and feedback — that leave observable signatures in their stellar light distributions, such as sizes, surface brightness profiles, and morphologies, over cosmic time. Quantifying these structural properties and tracing their evolution therefore provides a direct observational route to constraining the physical mechanisms

that govern galaxy formation and evolution (e.g., R. J. Buta 2013; C. J. Conselice 2014).

A long-established and widely adopted approach to characterizing galaxy structure is to decompose their light profiles into two primary components: a centrally concentrated bulge and an extended disk (P. D. Allen et al. 2006; L. Simard et al. 2011; C. N. Lackner & J. E. Gunn 2012; J. T. Mendel et al. 2014; M. Vika et al. 2014; R. Kennedy et al. 2016; R. Lange et al. 2016; P. Dimauro et al. 2018; C. Bottrell et al. 2019; B. Häußler et al. 2022; S. Casura et al. 2022; L. Quilley & V. de Lapparent 2022; L. Quilley & V. de Lapparent 2023; E. Merlin et al. 2023; Euclid Collaboration et al. 2023; K. Jegatheesan et al.

Email: angelogeorget@gmail.com

* Canada Research Chair

2024; L. Quilley et al. 2025; L. Kawinwanichakij et al. 2025, among many others). Empirically, disks are typically larger in size and well described by exponential light profiles, while bulges are more compact and centrally concentrated (K. C. Freeman 1970; J. Kormendy 1977; C. M. Carollo et al. 1997; P. Lang et al. 2014; J. Méndez-Abreu et al. 2021; A. S. G. Robotham et al. 2022a; C. Lima-Dias et al. 2024; L. Kawinwanichakij et al. 2025). Beyond their distinct spatial distributions, bulges and disks also differ systematically in their stellar population properties — such as ages, star formation activity, and chemical enrichment — indicating that they trace different modes of mass assembly and evolutionary pathways (e.g., J. Kormendy & R. C. Kennicutt 2004; S. Khochfar & J. Silk 2006; D. B. Fisher & N. Drory 2011; C. Tonini et al. 2016; B. Clauwens et al. 2018; L. Zhang et al. 2025).

An important and widely used scaling relation for linking galaxy structure to stellar mass assembly is the size–mass relation (SMR). Because it connects mass growth to structural growth, the SMR provides a sensitive probe of the physical processes — such as mergers, tidal interactions, emergence of bulges, fading of disks, etc. — that shape galaxies over cosmic time. It is now well established that the total sizes of galaxies scale with their stellar mass, and that star-forming galaxies (SFGs) and quiescent galaxies (QGs) have their own sequences in the size–mass plane (e.g., S. Shen et al. 2003; I. Trujillo et al. 2004; Y. Guo et al. 2009; R. J. Williams et al. 2010; A. B. Newman et al. 2012; A. van der Wel et al. 2014; R. Lange et al. 2015; A. L. Faisst et al. 2017; N. Roy et al. 2018; L. A. Mowla et al. 2019; J. Matharu et al. 2019, 2020, 2024; I. Damjanov et al. 2019, 2023; L. Kawinwanichakij et al. 2021, 2025; S. E. Cutler et al. 2022; A. George et al. 2024, 2025; K. Ito et al. 2024; M. Martorano et al. 2024; E. Ward et al. 2024; K. Ormerod et al. 2024).

With the advent of deep, high-resolution imaging, it has become possible to extend this analysis to individual structural components. Recent studies have demonstrated that both bulges and disks follow their own SMRs, which share a similar linear functional form in logarithmic space but differ in slope and normalization (V. A. Bruce et al. 2014; P. Dimauro et al. 2019; A. Hashemizadeh 2022; D. S. H. Hon et al. 2023; K. V. Nedkova et al. 2024; R. H. W. Cook et al. 2025; L. Kawinwanichakij et al. 2025, among others). This implies that the observed SMRs of whole SFGs and QGs largely reflect the combined contributions of their bulge and disk components. Component-resolved SMRs therefore offer a powerful framework for testing theoretical models of galaxy formation, as they place direct constraints on

how stellar mass is assembled spatially within galaxies and how different physical processes shape bulges and disks across cosmic time.

Despite significant progress, existing studies of bulge and disk SMRs remain limited by sample size, survey area, and redshift coverage. Many of the component-resolved analyses are based on a few hundred to a few thousand galaxies, typically drawn from deep but narrow fields ($\lesssim 2 \text{ deg}^2$), which makes them susceptible to cosmic variance and limits statistical power (e.g., V. A. Bruce et al. 2014; P. Dimauro et al. 2019; J. Méndez-Abreu et al. 2021; A. Hashemizadeh 2022; D. S. H. Hon et al. 2023; K. V. Nedkova et al. 2024). In addition, many studies probe only a small restricted redshift interval or focus on single epochs, limiting their ability to robustly trace the evolution of bulges and disks (e.g., J. Méndez-Abreu et al. 2021; A. S. G. Robotham et al. 2022a; D. S. H. Hon et al. 2023). These limitations motivate the need for wide-area surveys that enable component-resolved structural measurements for large, representative samples out to intermediate redshifts.

So far, works in the literature have primarily looked at the SMR of bulges and disks only in a single rest-frame wavelength. The preferred choice so far has been in the optical/near-IR regime, as the light in these wavelengths traces the distribution of old and low-mass stars that form the bulk of the stellar population in galaxies (G. Bruzual & S. Charlot 2003; D. C. Martin et al. 2005). However, in our pilot study (A. George et al. 2024, hereafter Paper I), we showed that it is also vital to analyze galaxy morphologies in rest-frame UV to study the distribution of young and/or low-metallicity stars. By contrasting the structural properties of the rest-UV and rest-optical, we could identify the physical processes driving the galaxy evolution. For example, the growth driven by minor mergers should increase the average of sizes of galaxies with similar mass faster in rest-UV than in the rest-optical as shown by our earlier works, Paper I and A. George et al. (2025, hereafter Paper II). Similarly, inside-out quenching — due to AGN feedback, growth of bulges, etc. — should make the galaxy sizes smaller in the longer wavelength.

Therefore, in this work, for the first time, we analyze the morphologies and the redshift evolution of the SMR for bulges and disks of SFGs and QGs in two rest-frame wavelengths (UV and optical) over 6 billion years ($0.1 < z < 0.85$). We utilize the deep multi-wavelength dataset from the CFHT CLAUDS (M. Sawicki et al. 2019) and Subaru HSC-SSP (H. Aihara et al. 2018, 2019) surveys that cover a sky area of 18.6 deg^2 .

Quantifying the structural evolution of bulges and discs requires a choice of mass-anchor, with existing

studies generally adopting one of two complementary strategies. Some studies compute component SMRs using bulge and disk stellar masses (e.g., V. A. Bruce et al. 2014; P. Dimauro et al. 2019; A. Hashemizadeh 2022; A. S. G. Robotham et al. 2022a; K. V. Nedkova et al. 2024; R. H. W. Cook et al. 2025), while others relate component sizes to the total stellar mass of the host galaxy (e.g., C. Tonini et al. 2016; J. Méndez-Abreu et al. 2021; L. Kawinwanichakij et al. 2025). These complementary approaches respectively probe the intrinsic evolution of individual components and their growth in the context of the galaxy as a whole.

To consistently compare structural components across this large and diverse sample, we adopt the total stellar mass of the galaxy, rather than the component mass, when analyzing the SMRs of bulges and disks. Since stellar mass is a fundamental property closely linked to various galaxy characteristics (E. F. Bell et al. 2003; B. Panter et al. 2004; I. K. Baldry et al. 2006; Y.-j. Peng et al. 2010; S. Zibetti & A. R. Gallazzi 2022), this approach allows us to compare the characteristic sizes of bulges and disks in galaxies of similar stellar mass, and, thus, similar global properties.

This paper is divided into the following sections. We provide an overview of the dataset used in Section 2. Section 3 describes the methodology we adopt to measure the morphological properties (especially the size) of the structural components of the SFGs and QGs through decomposition into a bulge and disk in two rest-frame wavelengths. We then analyze the size-stellar mass relation for bulges and disks in Section 5 and the evolution of the characteristic sizes (SMR zero points) for these structural components in Section 6. Section 7 provides the interpretation of the results for SFGs (Section 7.1) and QGs (Section 7.2). We address potential caveats and methodological limitations in Section 8. Finally, we summarize our major findings in Section 9.

Throughout this work, we use M to denote stellar mass in solar units (M_*/M_\odot) unless otherwise stated. Additionally, we use the AB magnitude system and adopt a standard cosmology with $\Omega_M = 0.3$, $\Omega_\Lambda = 0.7$ and $H_0 = 70$ km/s/Mpc throughout this paper.

2. DATA

2.1. Dataset Overview and Galaxy Classification

For this work, we utilize the same data images and ancillary data as in Paper II. In short, we use the Deep and UltraDeep multi-band ($U + griz$) imaging data from the CFHT Large Area U -band Deep Survey⁸ (CLAUDS; M.

Sawicki et al. 2019) and the Hyper Suprime-Cam Subaru Strategic Program⁹ (HSC-SSP; H. Aihara et al. 2018, 2022) covering 18.6 deg² spanning over four survey fields (E-COSMOS, DEEP2-3, ELAIS N1 and XMM-LSS). We also have complementary imaging from the SHIRAZ Survey (M. Annunziatella et al. 2023, IRAC 3.6m and 4.5m) in the E-COSMOS, DEEP2-3 and ELAIS N1 fields, and the VIRCAM near-IR wavelength coverage (Y , J , H , and K_s bands) from the VIDEO (M. J. Jarvis et al. 2013) and UltraVISTA (H. J. McCracken et al. 2012) surveys over a combined 5.5 deg² across the E-COSMOS and XMM-LSS fields.

Similar to Paper II, we rely on the CLAUDS+HSC photometric catalog by G. Desprez et al. (2023) made with the HSC Pipeline (hscPipe; J. Bosch et al. 2018) to select our galaxy sample. Photometric redshift estimates are based on $Ugrizy$ or $Ugrizy+JHK_s$ (where VISTA JHK_s bands are available) and calculated using Phosphoros¹⁰ (Euclid Collaboration et al. 2020; Paltani et al. in preparation), a template fitting code created to be part of the Euclid photo- z pipeline. These photo- z measurements have a precision of 0.05 and a very small fraction of outliers ($< 6\%$ at $i < 25$; G. Desprez et al. 2023).

Estimation of other galaxy properties that we use (stellar masses and rest-frame colors) are based on $Ugrizy$ photometry together with VISTA JHK_s (in the XMM-LSS field) or IRAC Ch1 and Ch2 (in the other three fields) using LePhare (O. Ilbert et al. 2006) template-fitting code. This is done by fixing the redshifts of galaxies to the photometric redshift values from G. Desprez et al. (2023) and assuming the Chabrier initial mass function (IMF; G. Chabrier 2003) and an exponentially decaying star formation history. Following the approach of T. Moutard et al. (2016), dust extinction is modeled using a range of attenuation laws: M. L. Prevot et al. (1984), D. Calzetti et al. (2000), and a modified Calzetti curve that incorporates a bump at 2,175Å. We will present the detailed procedure for these estimations in L. Chen et al. (in preparation).

As in Paper II, we adopt a modified variant of the $NUVrK$ color-color selection (S. Arnouts et al. 2013; T. Moutard et al. 2018, 2020) to classify our galaxies into star-forming and quiescent (L. Chen et al., in preparation). In short, to smoothly account for redshift and mass dependence, we adopt an $NUVrK$ selection method based on the distribution of galaxies in a four-dimensional parameter space defined by $NUV-r$, $r-K$,

⁹ <https://hsc.mtk.nao.ac.jp/ssp/>

¹⁰ <https://anaconda.org/astrorama/phosphoros>

⁸ <https://www.clauds.net/>

254 redshift and stellar mass. Our selection criteria for QGs
255 are

$$256 \quad NUV - r > (c_1 z + c_2 \log M + c_3)[(r - K) + c_4], \quad (1)$$

$$256 \quad NUV - r > c_5,$$

257 where the empirically calibrated values for constants
258 $c_1 - c_5$ are 3.868, 2.397, 1.003, -0.278 and 0.039, respec-
259 tively. Full details of the SFG/QG classification method
260 for the CLAUDS+HSC sample will be available in L.
261 Chen et al. (in preparation).

262 2.2. Sample Mass Completeness

263 We estimate the stellar mass completeness using the
264 method of L. Pozzetti et al. (2010), which infers the
265 limiting stellar mass by rescaling the stellar masses of
266 individual galaxies to the survey flux limit and defining
267 the mass above which a fixed fraction (85%) of galaxies
268 would be observable. The resulting completeness limits
269 are shown in Figure 1, with SFGs on the left and QGs
270 on the right. The parent catalog is flux-limited to $i <$
271 26 to ensure reliable photo- z and contains stellar mass
272 estimates for these galaxies. Based on this limit, the
273 catalog is mass-complete down to $\log M \gtrsim 8.5$ for SFGs
274 and $\log M \gtrsim 9$ for QGs at $z \leq 0.85$. These completeness
275 limits are indicated by the yellow curves in the figure.

276 However, bulge-disk decomposition is not feasible for
277 all galaxies in the parent catalog. As we show using
278 simulations — similar to those in our previous works
279 (Papers I & II) and discussed in Appendix A — we can
280 reliably model structural components for bright galaxies
281 with $i < 22$ mag. We thus restrict our structural analy-
282 sis to this brighter subset. This i -band cut increases the
283 stellar mass completeness thresholds to $\log M \gtrsim 10.8$
284 for SFGs and $\log M \gtrsim 11.0$ for QGs at $z > 0.7$. The
285 completeness limits of the final SFG and QG samples
286 we use throughout this study are represented by green
287 curves in Figure 1.

288 Adopting a single mass limit across the full red-
289 shift range would require using the highest complete-
290 ness threshold, corresponding to $\log M \gtrsim 10.8$ for SFGs
291 and $\log M \gtrsim 11.0$ for QGs, which would severely re-
292 duce the available sample size. To retain a statistically
293 large sample while maintaining mass completeness, we
294 therefore adopt redshift-dependent stellar mass limits,
295 defined within discrete redshift bins motivated by the
296 SMR analysis in this work and in Paper II.

297 The black boxed regions in Figure 1 represent the se-
298 lected sample with these evolving stellar mass limits.
299 For SFGs, the adopted mass limits are: $\log M > 9.7$
300 ($0.1 < z < 0.4$), > 10.2 ($0.4 < z < 0.55$), > 10.5
301 ($0.55 < z < 0.7$), and > 10.8 ($0.7 < z < 0.85$).
302 For QGs, the corresponding limits are: $\log M > 10.4$

Category	Number
All objects ($i < 26$)	7,493,735
Galaxies with reliable mass measurements	5,239,677
$i < 23$	873,867
SE++ fits	569,464
UV	547,906
Optical	531,883
$i < 22$	262,168
$0.1 < z < 0.85$	221,575
$\log M > 9.7$	139,809
SFGs	84,420
QGs	55,389
Mass-complete sample:	
SFGs	44,086
QGs	34,108
Reliable SE++ fits in UV:	
SFGs	42,245
SFG Bulges [†]	16,654
SFG Disks [‡]	41,795
QGs	33,319
QG Bulges [†]	25,664
QG Disks [‡]	32,597
Reliable SE++ fits in Optical:	
SFGs	42,341
SFG Bulges [†]	21,376
SFG Disks [‡]	41,374
QGs	33,470
QG Bulges [†]	30,415
QG Disks [‡]	28,899
Reliable SE++ fits in UV & Optical:	
SFGs	41,210
QGs	32,885

Table 1. Sample selection from the CLAUDS+HSC data.

[†]Number of bulges with $B/T > 0.2$, used for SMR fits.

[‡]Number of disks with $B/T < 0.8$, used for SMR fits.

303 ($0.1 < z < 0.55$), > 10.7 ($0.55 < z < 0.7$), and > 11.0
304 ($0.7 < z < 0.85$). We limit our QG sample with stellar
305 mass $\log M > 10.4$ as we intend to analyze the QG SMR
306 above its pivot point ($\log M \sim 10.4$; Papers I & II). We
307 thus obtain a mass-complete sample of 44,086 SFGs and
308 34,108 QGs, as summarized in Table 1.

309 3. METHODOLOGY

310 To investigate the evolution of galaxy structural
311 components separately, we perform two-component
312 bulge+disk decompositions using SOURCEXTRACTOR++¹¹ (SE++; E. Bertin et al. 2020), a modern
313

¹¹ <https://ascl.net/2212.018>

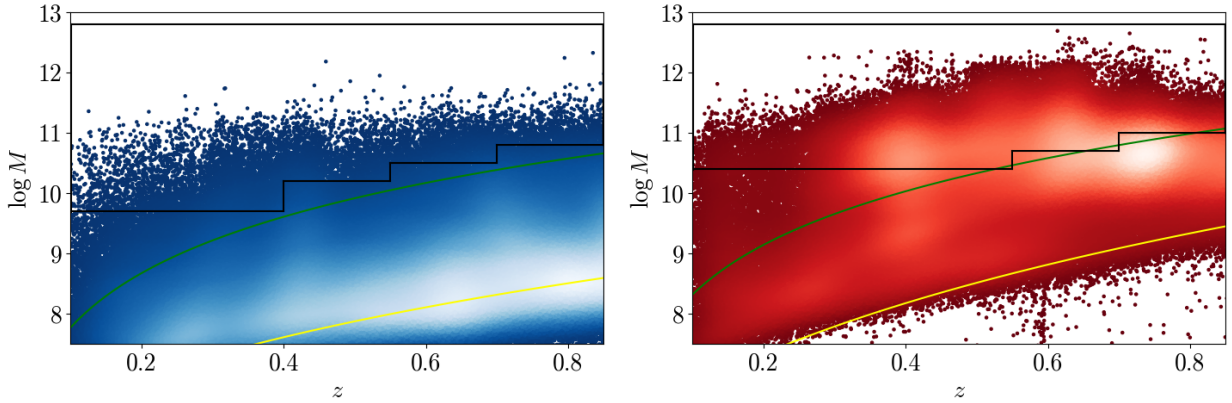


Figure 1. Left: Selection of mass complete sample for SFGs based on the magnitude cut ($i < 22$ mag) adopted in this work. The distribution of the SFGs from the parent catalog in the stellar mass-redshift space is shown in blue heatmap, and the yellow curve shows the 85% mass completeness limit for these SFGs. The 85% mass completeness limit after applying the magnitude cut is shown in green. The selected mass complete sample is shown within a black boxed region. Right: Selection of mass complete sample for similarly bright QGs, with mass completeness limit for QGs shown in green. The distribution of the QGs from the parent catalog is shown in red heatmap. Similar to the left panel, the selected QG sample is shown within the boxed region.

314 extension of SOURCE EXTRACTOR¹² (SEXTRACTOR; E.
 315 Bertin & S. Arnouts 1996; E. Bertin 2011), that enables
 316 simultaneous multi-component, multi-band structural
 317 fitting. Unlike SExtractor, which is primarily designed
 318 for source detection and basic photometry, SE++
 319 allows parametric modeling of galaxy light profiles
 320 across multiple wavelengths in a unified framework.
 321 Using SE++, we model each galaxy with a Sérsic
 322 profile for the bulge component, where the Sérsic
 323 index n is left free, and an exponential profile for the
 324 disk component, corresponding to a fixed Sérsic index
 325 $n = 1$. This bulge+disk fits are performed across the
 326 $Ugriz$ bands of the CLAUDS+HSC imaging separately.
 327 This multi-component fitting approach is analogous
 328 to our previous single-Sérsic analyses (Papers I & II),
 329 but enables us to isolate the sizes of bulges and disks
 330 individually and assess their distinct contributions to
 331 the observed structural evolution of galaxies.

332 Prior to structural fitting, SE++ identifies and seg-
 333 ments galaxies using a model-based detection and de-
 334 blending scheme. The software models light distribu-
 335 tions by comparing observed pixel values to a paramet-
 336 ric model convolved with local PSFs, which we generate
 337 for each field using the Point Spread Function Extrac-
 338 tor¹³ (PSFEX; E. Bertin 2011, 2013) and SEXTRACTOR.
 339 The fitting process minimizes a modified least-
 340 squares cost function weighted by pixel-level uncertain-
 341 ties. To ensure robustness against neighboring contam-
 342 ination or noise spikes, we set the modified χ^2 scale to

343 4; this processes residuals through a monotonic trans-
 344 formation that downweights large deviations to prevent
 345 outliers from biasing the fit. Furthermore, we utilize the
 346 iterative model-fitting mode with a maximum of 250 it-
 347 erations and 5 meta-iterations to ensure stable conver-
 348 gence in these complex bulge-disk decompositions.

349 The SE++ configuration is optimized for robust fit-
 350 ting of galaxies with diverse morphologies. In gen-
 351 eral, we adopt the SE++ configurations and method-
 352 ologies followed by the Euclid Collaboration recently
 353 (Euclid Collaboration et al. 2023; L. Quilley et al.
 354 2025). We adopt a detection threshold of 5σ above
 355 the local background, estimated on a grid with
 356 `background-cell-size = 256` pixels to balance accu-
 357 racy in crowded fields. The fitting model combines a
 358 Sérsic bulge ($x, y, \text{flux}_B, R_{e,B}, n_B, \text{pa}, q_B$) and
 359 an exponential disk ($x, y, \text{flux}_D, R_{e,D}, \text{pa}, q_D$),
 360 where R_e is the effective radius along the semi-major
 361 axis, n is Sérsic index, pa is the position angle and q
 362 is the axis ratio. We use initial parameters derived from
 363 built-in functions for isophotal centroid, flux, size and
 364 position angle.

365 Beyond its robust cost function, SE++ offers ad-
 366 vanced regularization capabilities that stabilize multi-
 367 component fits — a significant improvement over the
 368 original SEXTRACTOR. We leverage these features by
 369 defining physically motivated parameter ranges and in-
 370 formative Gaussian priors on the model parameters. For
 371 the bulge component, we implement a non-Gaussian
 372 prior on the Sérsic index, $N = \exp(n - n_0)$ with $n_0 = 6$.
 373 This serves to penalize unphysically high concentrations
 374 ($n > 7$) while remaining uninformative for lower indices
 375 ($n \lesssim 5$), a more conservative threshold than the $n_0 = 8$

¹² <https://ascl.net/1010.064>

¹³ <http://ascl.net/1301.001>

adopted in a recent single Sérsic SE++ analysis by [M. Pizzardo et al. \(2026\)](#). To maintain a structural hierarchy, we also apply a joint prior on the relative sizes of the components: $\Delta \log R_e = \log_{10}(R_{e,B}) - (2 \log_{10}(R_{e,D}) - 0.48)$. This regularization prevents non-physical solutions, such as component swapping, while allowing the final sizes to remain data-driven. We further stabilize the fit using a logit-normal prior on the bulge-to-total (B/T) ratio centered at $B/T \approx 0.1$ and constrain all axis ratios and radii using exponential ranges to avoid degenerate solutions.

Similar to Papers I & II, we test the robustness of the fitting methodology and estimate the uncertainties in the fit parameters obtained from SE++ through extensive simulations. We are able to model more than 90% of the simulated galaxies, and we use the spread in the measured parameters of these galaxies to estimate the uncertainties. Additionally, we also use the simulations to correct the biases present in our size measurements. Throughout this work, we adopt these corrected size estimates for analysis. We discuss these simulations, their results, uncertainty estimation and bias correction in [Appendix A](#).

4. BULGES AND DISKS IN TWO REST-FRAME WAVELENGTHS

Following the methodology prescribed in Papers I & II, we estimate the structural parameters of bulge and disk components in two rest-frame wavelengths: 3000Å (UV) and 5000Å (optical). In short, we take a weighted average of structural parameters of a galaxy in the two nearest observed bands that either bracket or are close to the rest-frame 3000Å and 5000Å. In case of fit results available in only a single band that is close to the rest-frame wavelength, we simply adopt the results in that band. Thus, the observed bands we use to estimate the Sérsic parameters of bulges and disks in the same rest-frame regimes depend on the redshift of individual galaxies.

We have modeled the light profiles of 569,464 galaxies with $i < 23$ mag in the CLAUDS+HSC images, although after the cuts based on magnitude ($i < 22$ mag), redshift ($0.1 < z < 0.85$) and stellar mass for completeness, the final sample with bulge+disk decomposition used for analysis in this work contains 44,086 SFGs and 34,108 QGs. However, we note that not all galaxies have light profiles modeled in both rest-frame wavelengths. Thus, we have structural parameters for 42,245 SFGs in rest-UV and 42,341 SFGs in rest-optical. Similarly, the corresponding numbers for QGs are 33,319 and 33,470, respectively. Out of these, 41,210 SFGs and 32,885 QGs have structural component fits in both rest-frame wave-

lengths. We provide these numbers in [Table 1](#). Because SE++ failed to model $\sim 5\%$ of the sample, excluding these “failed” galaxies won’t affect our analysis.

The bulge and disk sizes used in this analysis have been calibrated against systematic offsets identified in our image simulations ([Appendix A](#)). For example, the pipeline, in general, overestimates the bulge sizes. Although this bias reaches $\sim 50\%$ if the measured $n_B > 4.5$, it is significantly lower ($\sim 10\%$) for bulges with $n_B < 4.5$ where majority of the observed galaxies are ([Figure 11](#)). In contrast, the biases in the measured disk sizes are significantly smaller than those of bulge sizes except for bulge dominated galaxies ($B/T > 0.75$) that do not constitute a dominant population in our disk analysis sample ([Figure 12](#)). Accordingly, we apply a correction to our bulge and disk size measurements to account for these systematic biases, as detailed in [Appendix A.2](#). While all results presented hereafter utilize these bias-corrected measurements, we note that the qualitative trends of this work remain consistent even without these adjustments.

In addition, our simulations ([Appendix A](#)) indicate that although SE++ effectively models the overall galaxy light profile, the size measurements of individual structural components (i.e., bulge and disk) are reliable only when the component contributes at least 20% of the total modeled flux. In other words, we can robustly recover the size of a disk only if it is not overwhelmed by the bulge (and vice versa). Thus, in this work, we introduce cuts in the bulge-to-total (B/T) ratios while selecting galaxy bulges and disks for analysis. For bulge and disk components, these selection criteria are $B/T > 0.2$ and $B/T < 0.8$, respectively. The number of bulges and disks that meet these criteria is listed in [Table 1](#). Similar B/T cuts have been adopted in previous studies (e.g., [P. D. Allen et al. 2006](#); [A. Meert et al. 2015](#); [P. Dimauro et al. 2018, 2022](#); [B. Häußler et al. 2022](#); [Euclid Collaboration et al. 2023](#); [A. Ghosh et al. 2023](#)). Because the excluded bulges and disks have mass and redshift distributions similar to those of the parent bulge and disk populations, these B/T thresholds are not expected to introduce significant biases. Ultimately, our analysis focuses on the size evolution of bulges and disks in galaxies for which these components are sufficiently prominent to be reliably measured.

[Figure 2](#) presents the distribution of effective radii for these *significant* bulge and disk components of SFGs (the top panels) and QGs (the bottom panels), measured in two rest-frame wavelengths. In this figure and throughout this work, we use cyan (UV) and blue (optical) colors for SFG disks and pink (UV) and red (optical) colors for QG disks. Likewise, we use green (UV) and

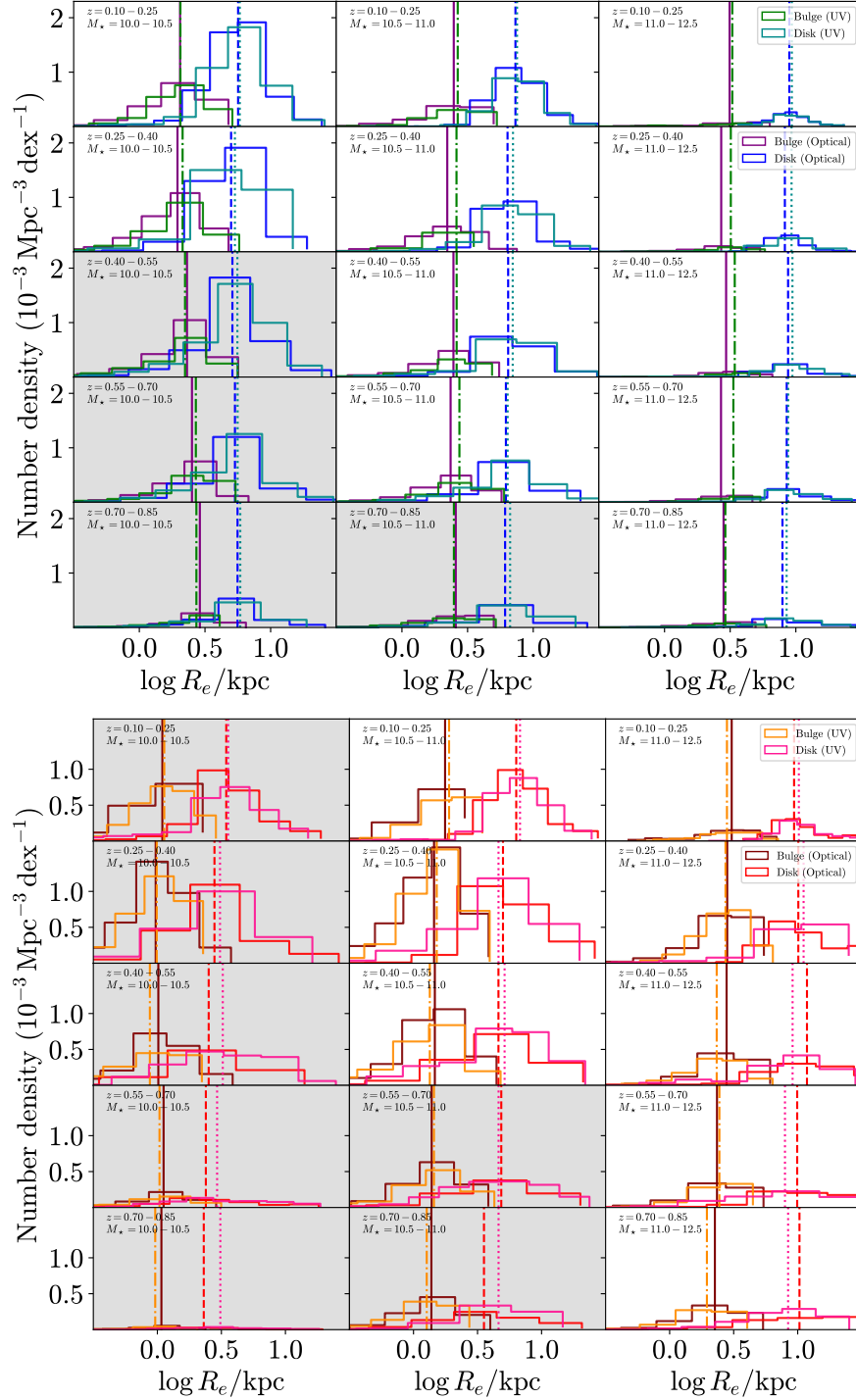


Figure 2. Top panels: Distribution of the sizes of the structural components of SFGs in rest-frame UV and optical wavelengths. The blue (purple) color denotes bulges in UV (optical), whereas the cyan (green) shows the distribution for disks in UV (optical). As we move from left to right, the stellar masses of the galaxies increase. Similarly, we move from lower to higher redshift bins downward. The y-axis shows the number density in comoving volume. The median values of $\log(R_e/\text{kpc})$ for bulges and disks in UV and optical wavelengths in every panel are denoted by the vertical lines. We show the sizes of bulges only if the bulge-to-total ratio, $B/T > 0.2$, and we show the sizes of disks only if $B/T < 0.8$. The panels affected by sample mass incompleteness are shaded grey. Bottom panels: Distribution of the sizes of the structural components of QGs. The red (maroon) color represent bulges in UV (optical), while the pink (orange) represents disks in UV (optical).

479 purple (optical) colors for SFG bulges, and orange (UV)
 480 and maroon (optical) colors for QG bulges.

481 As expected, disks are generally more extended than
 482 bulges across both galaxy types and wavelengths, as in-
 483 dicated by the differences in their median sizes (vertical
 484 lines in Figure 2). This trend holds across all stellar
 485 mass ($\log M_\star > 10$) and redshift ($0.1 < z < 0.85$) bins.
 486 Disk size distributions are also significantly broader than
 487 those of bulges. In many bins, a non-negligible frac-
 488 tion of disks have effective radii smaller than the me-
 489 dian bulge size, whereas bulges with sizes exceeding the
 490 median disk size in the same bin are extremely rare.

491 The sizes of bulges and disks depend on the stellar
 492 mass and redshift of galaxies. As expected, more mas-
 493 sive galaxies tend to have larger bulges and disks irre-
 494 spective of the wavelength or redshift bin. For example,
 495 in the rest-optical, SFG bulges (disks) at $0.4 < z < 0.55$
 496 are 28% (71%) larger in the highest mass bin than in
 497 the lowest mass bin (third row in the top panels in Fig-
 498 ure 2). For QGs, the differences are 174% and 370%
 499 for bulges and disks, respectively. The sizes of bulges
 500 and disks of both SFGs and QGs in the mass complete
 501 bins show a general negative trend with redshift in both
 502 wavelengths. As an example, the bulges and disks of
 503 the most massive ($\log M > 11$) SFGs grew by 12% and
 504 13%, respectively, between the highest and lowest red-
 505 shift bins in rest-optical (second column in the top pan-
 506 els in Figure 2). These trends with stellar mass and
 507 redshift reflect the similar trends we observe for global
 508 sizes of galaxies in our Paper I.

509 The sizes of galaxy components also depend on the
 510 rest-frame wavelength. Based on the median values in
 511 Figure 2, the disks are generally more extended in the
 512 rest-UV than in the optical wavelength, except for most
 513 massive QGs at $z > 0.4$. Bulges of SFGs tend to be larger
 514 in the shorter wavelengths in the mass-complete bins. In
 515 contrast, QGs tend to have larger bulges in the longer
 516 wavelength in many of the bins.

517 We caution, however, that the plotted distributions do
 518 not account for uncertainties in the size measurements.
 519 While they effectively highlight overall trends, a more
 520 rigorous interpretation requires incorporating both mea-
 521 surement uncertainties and systematic effects. We ad-
 522 dress this by using size estimates derived from the SMR,
 523 which accounts for uncertainties in both mass and size
 524 measurements when estimating the characteristic sizes
 525 of galaxy components at a given stellar mass (Section 5).

526 The distribution of Sérsic indices of bulges (n_B) differs
 527 between SFGs and QGs (the top panel of Figure 3). In
 528 general, SFG bulges exhibit lower Sérsic indices than QG
 529 bulges, with typical values of $n_B < 2$ for SFGs and $n_B >$
 530 2 for QGs. The SFG bulges tend to have more disk-like,

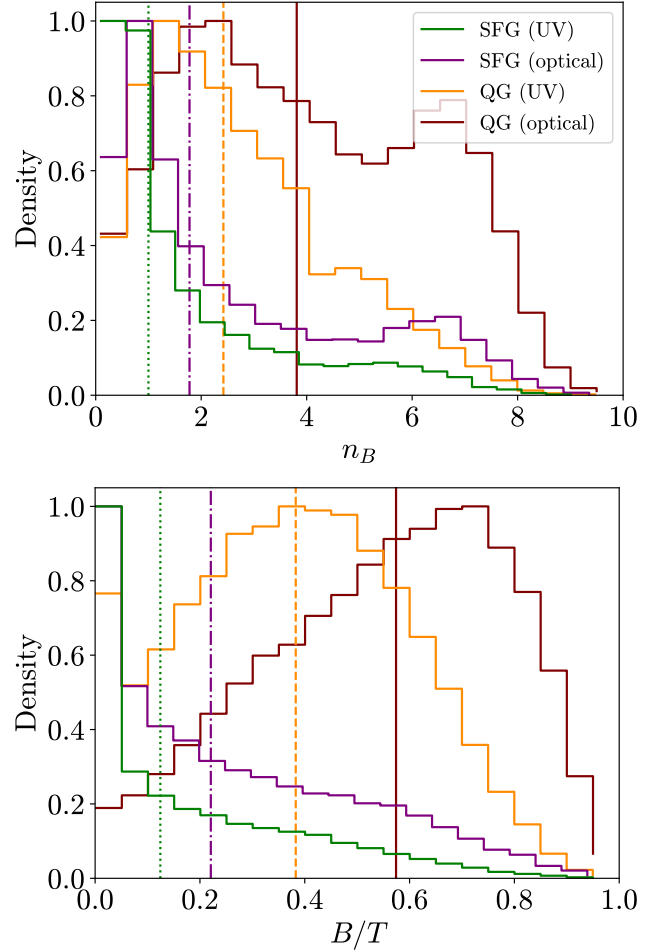


Figure 3. The top panel: Distribution of Sérsic indices of bulges in $\log M > 10.5$ galaxies with significant bulge component ($B/T > 0.2$). The green (purple) represent the SFG bulges in UV (optical) while the orange (maroon) represent the QG bulges in UV (optical). In each histogram, the bins are normalized to the value of the most populated bin. The vertical lines represent the median values. The bottom panel: Distribution of bulge-to-total (B/T) ratios for SFGs and QGs with stellar masses $\log M > 10.5$ in two rest-frame wavelengths. Rest of the details are same as in the top panel.

531 less centrally concentrated light profiles, while the QG
 532 bulges are generally more centrally concentrated.

533 The bulge Sérsic index also varies with rest-frame
 534 wavelength. Both SFGs and QGs show an increase in
 535 n_B at longer wavelengths, indicating that light profiles of
 536 the bulge components are more centrally concentrated in
 537 the rest-optical than in the rest-UV. Specifically, the me-
 538 dian Sérsic index for $\log M > 10.5$ SFG bulges increases
 539 from 1.0 in the rest-UV to 1.8 in the rest-optical, while
 540 for $\log M > 10.5$ QG bulges it increases from 2.4 to 3.8.
 541 This trend is consistent with our previous single Sérsic
 542 analysis (Paper I). Since those single Sérsic profiles are
 543 strongly weighted by the highest surface brightness pix-

els near galaxy centers, their n -trend should follow the trend in n_B for non-negligible bulge components. This wavelength dependence of the central concentration of bulge profiles indicates that the older stellar populations, traced by rest-optical light, are more centrally concentrated in bulges than the younger stars, which dominate the rest-UV emission.

Despite the generally lower concentration of UV light, both SFGs and QGs display a high- n_B tail in the rest-UV distribution. Conversely, some optical bulges have $n_B \sim 1$. This suggests diversity in bulge properties, indicating that not all bulges are structurally similar. However, we caution that Sérsic index measurements derived from ground-based imaging are generally less robust than R_e estimates, particularly due to PSF effects (R. Davari et al. 2014; C. Dewsnap et al. 2023; A. George et al. 2024).

We also observe a secondary peak in the n_B distribution between $6 < n_B < 8$, particularly in the rest-optical. This feature likely reflects the influence of the non-Gaussian prior ($n_0 = 6$) applied during the structural fitting to prevent unphysically high concentrations (Section 3). In cases where the pixel data favors an extremely concentrated profile ($n \gtrsim 8$), the optimizer may settle within this penalized range. While this indicates that the highest n_B values should be interpreted as conservative lower limits, it does not affect our primary finding that QG bulges are significantly more concentrated than those of SFGs across both wavelengths.

The distribution of B/T ratios (the bottom panel of Figure 3) shows trends similar to those seen in the distributions of Sérsic indices of bulges. As expected, QGs exhibit higher B/T values than SFGs, indicating a more prominent bulge component in their light profiles. For example, 47% of SFGs with $\log M > 10.5$ have bulges contributing less than 20% of the total rest-optical light, compared to fewer than 10% among QGs.

Additionally, both SFGs and QGs tend to have higher B/T values in the rest-optical compared to the rest-UV, confirming that a higher fraction of older stellar population is located in bulges than in disks compared to the younger stellar population. For SFGs with stellar mass $\log M > 10.5$, the median B/T increases from 0.12 in the rest-UV to 0.22 in the rest-optical. In similarly massive QGs, the median values are 0.38 and 0.57, respectively. In terms of bulge dominance, 61% of QGs are bulge-dominated ($B/T > 0.5$) in the rest-optical, compared to 29% in the rest-UV. For SFGs, the corresponding fractions are significantly lower — 20% in the rest-optical and 11% in the rest-UV. However, as with Sérsic index measurements, we caution that B/T estimates are generally less robust than measurements of R_e of indi-

vidual components that contribute significantly to the total light profile.

5. SIZE–MASS RELATIONS OF BULGES AND DISKS

With the bulge and disk sizes measured in two rest-frame wavelengths, we now examine the size–mass relation (SMR) of these structural components separately. The SMR is well established for star-forming and quiescent galaxies (e.g., S. Shen et al. 2003; R. J. Williams et al. 2010; A. van der Wel et al. 2014; L. A. Mowla et al. 2019; L. Kawinwanichakij et al. 2021, 2025; A. George et al. 2024, 2025), and has also been reported for bulges and disks in previous studies (e.g., V. A. Bruce et al. 2014; Y. Sofue 2016; P. Dimauro et al. 2019; A. Hashemizadeh 2022; D. S. H. Hon et al. 2023; K. V. Nedkova et al. 2024; R. H. W. Cook et al. 2025; L. Kawinwanichakij et al. 2025). Here, leveraging a statistically large and homogeneous sample over $0.1 < z < 0.85$, we extend these studies by providing a unified, multi-wavelength analysis of bulge and disk SMRs, which has so far been limited by sample size, redshift range, or the lack of consistent rest-frame wavelength comparisons.

We describe the SMR generally as

$$\log R_e = \log R_0 + \alpha \log \left(\frac{M}{M_0} \right) + \sigma_{\log R_e}, \quad (2)$$

where α is the slope, R_0 is the characteristic size at a fiducial stellar mass M_0 and $\sigma_{\log R_e}$ is the intrinsic scatter in the linear regression. In this work, we use the effective radius (R_e) of the bulge and disk components and the total stellar mass (M) of the galaxy to construct the SMR. Similarly to our previous works, we adopt $M_0 = 5 \times 10^{10} M_\odot$. While the slope of the SMR for QGs exhibits a transition at a pivot mass of $\log M = 10 - 10.4$ (e.g., A. van der Wel et al. 2014; L. Kawinwanichakij et al. 2021; A. George et al. 2024, 2025), our sample consists exclusively of QGs above this threshold (see Section 2).

Following Paper II, we model the SMR of massive QGs by fitting a linear relation (Eq. 2) using the Bayesian regression tool LINMIX.¹⁴ This package employs a Markov Chain Monte Carlo approach to estimate the posterior distributions of the best-fitting parameters, R_0 and α . An advantage of LINMIX is that it accounts for measurement uncertainties in both stellar mass and effective radius during the fitting process.

In this work, for the first time, we investigate the SMRs for bulge and disk components of SFGs and QGs

¹⁴ <https://linmix.readthedocs.io/en/latest/src/linmix.html>

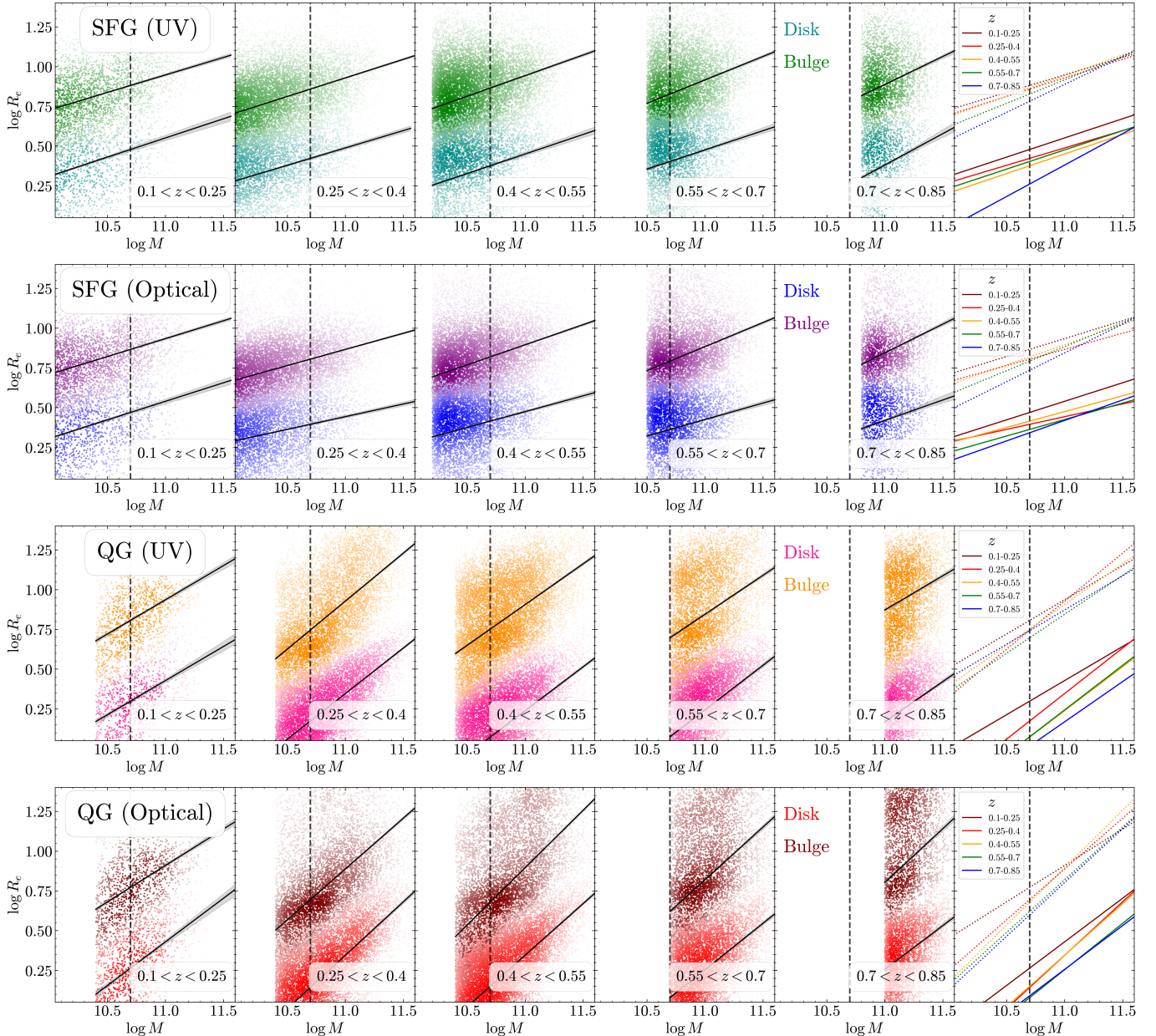


Figure 4. The SMRs for bulges and disks of SFGs (two rows on top) and QGs (two rows on the bottom). For SFGs and QGs, the first row shows the SMR in rest-UV and the second row shows the same in rest-optical. In each row, the first five panels show the SMRs in different redshift bins with the best-fit SMRs for the components are shown in black and uncertainty in gray band. In each redshift bin, the distribution of bulges and disks in the SMR plane is shown in separate colors as labeled in the image and we follow the same color pattern throughout this work. The final column illustrates the best-fit SMRs from all redshift bins together (solid lines for disks and dotted lines for bulges). In every panel, the vertical dashed line denotes the fiducial mass $M_0 = 5 \times 10^{10} M_\odot$.

642 separately in the rest-UV and rest-optical wavelengths. 649
 643 Based on our simulations, we restrict the SMR analysis 650
 644 of structural components to bulges with $B/T > 0.2$ 651
 645 and disks with $B/T < 0.8$. This is because unless a 652
 646 component contributes at least 20% of the total light of 653
 647 a galaxy, we are unable to reliably fit its Sérsic profile 654
 648 (Section 4).

We use the total stellar masses of galaxies (instead 655
 of their component stellar masses) to fit the SMRs for 656
 galaxy bulges and disks. While a multi-wavelength de- 657
 composition in the IR would ideally allow for individ- 658
 ual component mass estimates, such an approach is pre- 659
 cluded by the limited angular resolution of the IR data, 660
 which is up to three times coarser than that of the 661

Table 2. Best-fit parameters of the SMR (Eq. 2) and the size evolution (Eq. 3) of the bulge and disk components of SFGs and QGs in rest-frame UV and optical wavelengths.

Type	Comp.	z	$\log R_0$	α	$\sigma_{\log R_e}$	$\log M_t$	N_t	N_s	z_{med}	R_e^0	β
3000Å											
SFG	Bulge	0.10 – 0.25	0.48 ± 0.01	0.24 ± 0.02	0.16 ± 0.00	9.7	4671	1909	0.21	3.68 ± 0.67	−1.09 ± 0.45
		0.25 – 0.40	0.42 ± 0.01	0.22 ± 0.01	0.20 ± 0.00	9.7	15196	6098	0.31		
		0.40 – 0.55	0.38 ± 0.01	0.25 ± 0.02	0.22 ± 0.00	10.2	10971	3750	0.44		
		0.55 – 0.70	0.40 ± 0.01	0.24 ± 0.02	0.21 ± 0.00	10.5	7947	3471	0.63		
		0.70 – 0.85	0.26 ± 0.02	0.39 ± 0.05	0.23 ± 0.00	10.8	3460	1426	0.74		
	Disk	0.10 – 0.25	0.88 ± 0.01	0.22 ± 0.01	0.17 ± 0.00	9.7	4671	4581	0.20	8.50 ± 0.64	−0.53 ± 0.19
		0.25 – 0.40	0.86 ± 0.00	0.23 ± 0.00	0.17 ± 0.00	9.7	15196	15009	0.31		
		0.40 – 0.55	0.87 ± 0.00	0.26 ± 0.01	0.19 ± 0.00	10.2	10971	10878	0.45		
		0.55 – 0.70	0.83 ± 0.00	0.29 ± 0.01	0.19 ± 0.00	10.5	7947	7883	0.63		
		0.70 – 0.85	0.78 ± 0.01	0.35 ± 0.02	0.19 ± 0.00	10.8	3460	3444	0.74		
QG	Bulge	0.10 – 0.25	0.30 ± 0.01	0.43 ± 0.04	0.14 ± 0.00	10.4	1403	990	0.19	2.40 ± 0.35	−1.51 ± 0.35
		0.25 – 0.40	0.17 ± 0.00	0.57 ± 0.01	0.14 ± 0.00	10.4	9092	7630	0.35		
		0.40 – 0.55	0.07 ± 0.00	0.55 ± 0.01	0.19 ± 0.00	10.4	11485	8295	0.46		
		0.55 – 0.70	0.07 ± 0.01	0.56 ± 0.02	0.27 ± 0.00	10.7	7146	5605	0.65		
		0.70 – 0.85	0.02 ± 0.02	0.50 ± 0.03	0.22 ± 0.00	11.0	4193	3144	0.76		
	Disk	0.10 – 0.25	0.81 ± 0.01	0.43 ± 0.03	0.16 ± 0.00	10.4	1403	1346	0.18	6.76 ± 0.79	−0.49 ± 0.29
		0.25 – 0.40	0.75 ± 0.00	0.60 ± 0.01	0.17 ± 0.00	10.4	9092	8822	0.35		
		0.40 – 0.55	0.75 ± 0.00	0.51 ± 0.01	0.21 ± 0.00	10.4	11485	11308	0.46		
		0.55 – 0.70	0.70 ± 0.01	0.49 ± 0.02	0.27 ± 0.00	10.7	7146	6945	0.64		
		0.70 – 0.85	0.74 ± 0.02	0.43 ± 0.04	0.24 ± 0.00	11.0	4193	4176	0.76		
5000Å											
SFG	Bulge	0.10 – 0.25	0.47 ± 0.01	0.23 ± 0.02	0.18 ± 0.00	9.7	5220	2305	0.20	3.24 ± 0.34	−0.70 ± 0.26
		0.25 – 0.40	0.40 ± 0.01	0.16 ± 0.01	0.19 ± 0.00	9.7	15183	7348	0.31		
		0.40 – 0.55	0.41 ± 0.01	0.20 ± 0.02	0.19 ± 0.00	10.2	10923	5443	0.45		
		0.55 – 0.70	0.36 ± 0.00	0.21 ± 0.02	0.21 ± 0.00	10.5	7839	4438	0.63		
		0.70 – 0.85	0.34 ± 0.02	0.26 ± 0.05	0.18 ± 0.00	10.8	3176	1842	0.74		
	Disk	0.10 – 0.25	0.87 ± 0.00	0.22 ± 0.01	0.18 ± 0.00	9.7	5220	5091	0.20	8.08 ± 0.80	−0.64 ± 0.25
		0.25 – 0.40	0.80 ± 0.00	0.20 ± 0.00	0.16 ± 0.00	9.7	15183	14817	0.31		
		0.40 – 0.55	0.82 ± 0.00	0.25 ± 0.01	0.16 ± 0.00	10.2	10923	10704	0.45		
		0.55 – 0.70	0.79 ± 0.00	0.30 ± 0.01	0.17 ± 0.00	10.5	7839	7661	0.63		
		0.70 – 0.85	0.73 ± 0.01	0.36 ± 0.02	0.16 ± 0.00	10.8	3176	3101	0.74		
QG	Bulge	0.10 – 0.25	0.26 ± 0.01	0.54 ± 0.04	0.22 ± 0.00	10.4	1622	1268	0.18	2.07 ± 0.29	−1.02 ± 0.34
		0.25 – 0.40	0.14 ± 0.00	0.67 ± 0.01	0.24 ± 0.00	10.4	9131	8502	0.35		
		0.40 – 0.55	0.15 ± 0.00	0.65 ± 0.01	0.23 ± 0.00	10.4	11577	10442	0.46		
		0.55 – 0.70	0.08 ± 0.01	0.58 ± 0.02	0.26 ± 0.00	10.7	7196	6572	0.64		
		0.70 – 0.85	0.09 ± 0.02	0.55 ± 0.04	0.22 ± 0.00	11.0	3944	3631	0.76		
	Disk	0.10 – 0.25	0.77 ± 0.01	0.46 ± 0.02	0.16 ± 0.00	10.4	1622	1455	0.18	6.87 ± 0.39	−0.99 ± 0.14
		0.25 – 0.40	0.69 ± 0.00	0.64 ± 0.01	0.19 ± 0.00	10.4	9131	7414	0.35		
		0.40 – 0.55	0.68 ± 0.00	0.72 ± 0.01	0.23 ± 0.00	10.4	11577	10247	0.46		
		0.55 – 0.70	0.62 ± 0.01	0.66 ± 0.02	0.30 ± 0.00	10.7	7196	6273	0.64		
		0.70 – 0.85	0.60 ± 0.03	0.68 ± 0.06	0.31 ± 0.00	11.0	3944	3510	0.76		

NOTE—The column $\log M_t$ refers to the mass completeness cut in $\log M$. N_t refers to the total number of galaxies in the mass-complete sample with successful SE++ fits in a given redshift bin, while N_s is the number of galaxies used for the SMR fits ($B/T > 0.2$ for bulges and $B/T < 0.8$ for disks). The uncertainties below 0.005 are simply shown as 0.00. The uncertainties in z_{med} (median z) are all below 0.005.

656 CLAUDS+HSC imaging. Furthermore, structural pa- 655
 657 rameters vary significantly with wavelength (Section 4). 656
 658 Hence, by adopting total stellar mass, we avoid assigning 657
 659 parameters from the optical light profile decomposition 658
 660 to the (unresolved) IR components of the spectral en- 659
 661 ergy distribution (SED) fitting. This approach ensures 660
 662 a robust and uniform baseline for our analysis across 661
 663 the entire redshift range. Although this differs from the 662
 664 component-mass approach utilized in some of the simi-

655 lar works in the literature (e.g., A. Hashemizadeh et al. 656
 657 2022; A. S. G. Robotham et al. 2022a), it enables a more 658
 659 direct investigation into how the internal components of 659
 660 galaxies with similar stellar mass have evolved in size 660
 661 since $z \sim 0.85$. 661

662 We present the best-fit SMRs for the bulge and disk 662
 663 components of SFGs and QGs in Figure 4, with the cor- 663
 664 responding best-fit parameters provided in Table 2. In 664
 each row, the first five panels display the SMRs across

674 increasingly higher redshift bins, while the final panel
675 overlays the SMRs from all bins to facilitate comparison.
676 The top two rows show the component SMRs for SFGs
677 in the rest-frame UV and optical, respectively, while the
678 bottom two rows provide the same for the QG popula-
679 tion.

680 For both SFGs and QGs, bulges are consistently more
681 compact than their companion disks across both rest-
682 frame wavelengths. This size offset is more pronounced
683 in the rest-UV than in the optical, driven by the more
684 extended nature of disks and the increased compactness
685 of bulges at shorter wavelengths. In addition to redshift
686 evolution in the SMR slopes, the normalization at the
687 fiducial mass (M_0 is indicated by the vertical dashed
688 line) decreases at higher redshifts. This confirms that,
689 at fixed total stellar mass, both galaxy components were
690 significantly more compact in the past.

691 Figure 5 shows the best-fit SMR parameters (filled
692 symbols) and their evolution with redshift for both SFGs
693 and QGs. As expected, at a given fiducial stellar mass,
694 disks in both SFGs and QGs have larger characteristic
695 sizes (i.e., higher SMR zero points) than bulges, in both
696 rest-frame UV and optical wavelengths (top row of Fig-
697 ure 5). This size difference between disks and bulges is
698 more pronounced in the bluer band. For example, at
699 $0.4 < z < 0.55$, the disk of a fiducial SFG (QG) is 1.57
700 (2.39) times larger than the corresponding bulge in the
701 rest-frame optical. In the rest-frame UV, this difference
702 increases: the disk is 2.09 (3.79) times more extended
703 than the bulge in a similar SFG (QG).

704 Regardless of the wavelength, the sizes of bulges and
705 disks are larger for SFGs than for QGs. For example,
706 the bulges (disks) of $5 \times 10^{10} M_\odot$ SFGs are 62 – 91%
707 (26 – 48%) larger than those of QGs with the same
708 fiducial mass in all redshift bins at 5000\AA . P. Lang
709 et al. (2014) also report that the bulges and disks in
710 $\log M \sim 10.5$ SFGs are more extended than those in
711 similar QGs, although this offset is a factor of ~ 1.5
712 in a much broader redshift interval ($0.5 < z < 2.5$).

713 We also compare the evolution in the parameters of
714 the SMR for galaxy bulge and disk components to the
715 equivalent evolution of the SMR parameters for galaxies
716 fitted with a single Sérsic index (open circles in Figure 5;
717 Paper I). The zero points of galaxy disks are larger than
718 the SMR zero points from the single Sérsic component
719 measurements. In contrast, bulges are much more com-
720 pact than the total size of their host galaxies, in line with
721 other studies (e.g., L. Costantin et al. 2021). However,
722 for SFGs, the total galaxy size (from single-component
723 fit) closely resembles those of disk sizes, although the
724 offset between them increases with cosmic time. On the
725 other hand, the sizes of bulges and disks for QGs are

726 very different from the single-component size measure-
727 ments.

728 The top row in Figure 5 show a clear evidence for evo-
729 lution in the zero point of the SMR — i.e., the charac-
730 teristic size of bulges and disks at a fixed fiducial stellar
731 mass. Both bulges and disks exhibit size growth over
732 cosmic time. Between the redshift bins $0.7 < z < 0.85$
733 and $0.1 < z < 0.25$, the bulges of fiducial SFGs (QGs)
734 increase in size by 66% (91%) in the rest-frame UV,
735 and by 35% (48%) in the rest-frame optical wavelength.
736 Similarly, the disks of fiducial SFGs (QGs) grow by 26%
737 (17%) in the UV, and by 38% (48%) in the optical
738 wavelength. These results highlight two key trends: (1)
739 bulge sizes evolve more rapidly than disk sizes in the
740 rest-frame UV light, and (2) the rate of size growth is
741 wavelength-dependent. We explore these trends further
742 in Section 6.

743 The SMR slopes are steeper for QGs than for SFGs in
744 both wavelengths (middle row in Figure 5). The slopes
745 range from 0.16 to 0.36 for SFGs, while they are notably
746 steeper (0.43 – 0.72) for QGs. P. Dimauro et al. (2019)
747 also report similar SMR slopes for SFG disks and QG
748 bulges at these redshifts. The slopes we find for the
749 SFG components also agree with K. V. Nedkova et al.
750 (2024), but these authors report steeper slopes for the
751 QG bulges and shallower slopes for the QG disks. We
752 note that both of these works (P. Dimauro et al. 2019;
753 K. V. Nedkova et al. 2024) adopt the component mass
754 (instead of the total mass) in the SMR analysis.

755 In addition, the SMR relation tend to steepen for SFG
756 bulges and disks with redshift, a trend not observed for
757 SFG SMRs from single/global Sérsic profiles in this red-
758 shift range (e.g., A. van der Wel et al. 2014; L. Kawin-
759 wanichakij et al. 2021; A. George et al. 2024). In con-
760 trast, if we exclude the lowest redshift bin where our
761 sample size is the lowest, we find a decreasing trend in
762 the slopes for QG bulges (in UV and optical) and disks
763 (in UV). This is again different from the global SMR
764 trends reported in the literature (A. van der Wel et al.
765 2014; L. A. Mowla et al. 2019; L. Kawinwanichakij et al.
766 2021; K. V. Nedkova et al. 2021; A. George et al. 2024,
767 2025).

768 The intrinsic scatter in the SMR, denoted by $\sigma_{\log R_e}$,
769 is generally lower for the structural components (bulges
770 and disks) than for the global SMR of SFGs (left panels
771 in the bottom row of Figure 5). In contrast, QGs show
772 a comparable level of scatter between their global SMR
773 and the SMRs of their individual components, except for
774 their bulges in rest-UV (right panels). Overall, SFG disk
775 sizes exhibit a tighter correlation with stellar mass than
776 their bulge sizes across both wavelengths, as reflected
777 by the lower intrinsic scatter in their SMR. In contrast,

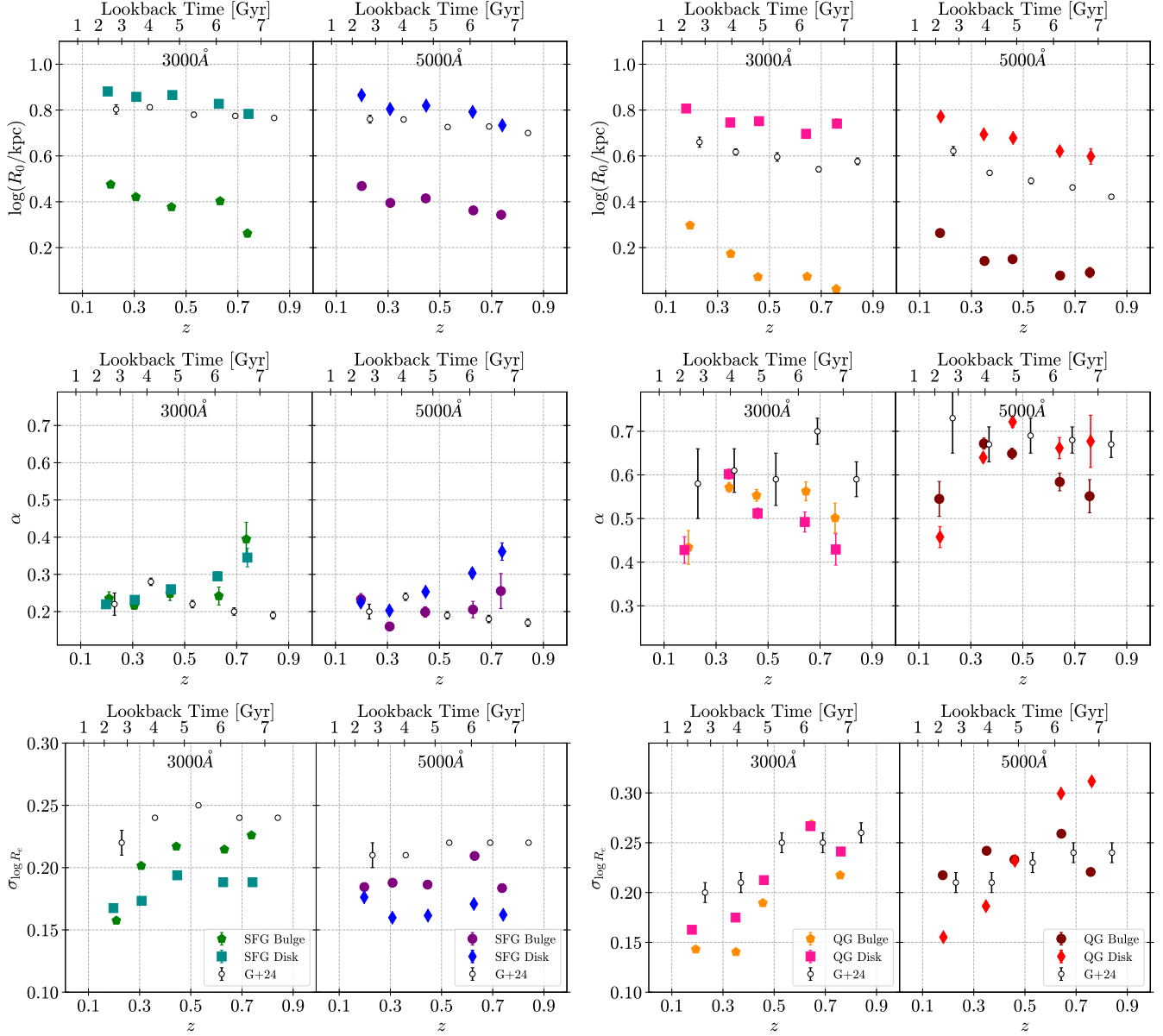


Figure 5. Redshift evolution of the best-fit SMR parameters for SFGs (left) and QGs (right). We plot the results in rest-UV and rest-optical for SFGs and QGs in separate sub-panels. Within each sub-panel, we show the SMR parameters for bulge and disk components separately. We also show the global SMR parameters based on the single Sérsic profile fitting from [A. George et al. \(2024, G+24\)](#) as open circles. Top row: Evolution of the SMR zero points ($\log R_0$). Middle row: Evolution of the SMR slopes ($\alpha = d \log R_e / d \log M$). Bottom row: Evolution of the intrinsic scatter in the SMR ($\sigma_{\log R_e}$).

778 QG bulges show smaller scatter than their disks in the
 779 rest-UV, whereas no such clear distinction exists in the
 780 rest-optical.

781 On average, the intrinsic scatter for SFG components
 782 ranges from 0.16 to 0.23, while for QG components it
 783 spans a broader range, from 0.14 to 0.31. For SFG
 784 disks, this scatter remains remarkably stable across the
 785 sampled redshift range, suggesting that disk growth is
 786 governed by well-regulated physical processes, such as
 787 continuous star formation and gas accretion. In con-

788 trast, the increasing scatter observed with redshift for
 789 SFG bulges (in rest-UV) and QG components (except
 790 for QG bulges in rest-optical) likely reflects a higher de-
 791 gree of structural diversity or more stochastic assembly
 792 histories at earlier epochs. The lack of evolution in the
 793 SMR scatter for optical QG bulges implies that the dis-
 794 tribution of old stellar population in them may be also
 795 driven by well regulated processes, just like SFG com-
 796 ponents. Finally, the total range of intrinsic scatter we
 797 observe for bulge and disk components (0.14 – 0.31) is

broadly consistent with the scatter found for the global SMRs of SFGs and QGs in Paper I, indicating that the stochasticity in component growth mirrors that of the galaxy as a whole.

In summary, our multi-wavelength analysis demonstrates that bulges and disks of SFGs and QGs follow distinct scaling relations that vary in slope, scatter, and wavelength-dependence. These systematic differences suggest that the internal components of both SFGs and QGs have been shaped by different physical assembly processes since $z \sim 0.85$. By establishing these robust relations at each epoch, we can now isolate the redshift evolution of component sizes at a fixed galaxy stellar mass to explore the physical mechanisms driving their growth.

6. SIZE EVOLUTION OF BULGES AND DISKS

We investigate the evolution of the structural components of galaxies as the Universe ages by tracking their characteristic sizes at a fiducial stellar mass of $5 \times 10^{10} M_{\odot}$, derived from their respective SMRs. We use SMR-based characteristic sizes because they are less biased than simple median sizes within stellar mass bins; the SMR fitting process explicitly accounts for measurement uncertainties in both size and stellar mass. We model this redshift evolution in two rest-frame wavelengths (3000Å and 5000Å) using a power-law model,

$$R_e = R_e^0(1+z)^\beta, \quad (3)$$

where R_e^0 is the characteristic size of a structural component of galaxies at $z = 0$.

As we describe in Section 5, the characteristic sizes of the bulge components are based on galaxies with $B/T > 0.2$ and the sizes of the disk components are based on galaxies with $B/T < 0.8$, to ensure the reliability of the component size measurements. Because of this, the number of galaxies used to fit the SMRs and, subsequently, the size evolution curve for the bulge and disk components are not the same. We provide the best-fit parameters in Table 2 and display the main results in Figure 6.

Consistent with our SMR analysis (Section 5), Figure 6 shows that disks (cyan/blue curves for SFGs and pink/red for QGs) are systematically more extended than the global galaxy size (black curves) across both rest-frame wavelengths, while bulges (green/purple curves for SFGs and orange/maroon for QGs) remain the most compact components.

Notably, we find that structural components grow at different rates than the galaxy as a whole (the left panel of Figure 6). For instance, the increasing offset between the sizes of SFG disks (blue/cyan) and global sizes

(black) over cosmic time in the left panel implies that disks expand more rapidly than the total galaxy profile ($\beta = -0.53 \pm 0.19$ vs $\beta = -0.26 \pm 0.13$ in UV; $\beta = -0.64 \pm 0.25$ vs $\beta = -0.35 \pm 0.14$ in optical). SFG bulges (green/purple) exhibit even steeper evolution, particularly in the rest-UV ($\beta = -1.09 \pm 0.45$ in UV; $\beta = -0.70 \pm 0.26$ in optical), suggesting they play a significant role in the structural transformation of star-forming systems. While these growth rates are subject to uncertainties, the trend of faster disk growth in the longer wavelength — contrasted with the opposite behavior in bulges — further underscores the distinct assembly histories of these two components.

In contrast, the evolutionary tracks for QG disks (pink/red curves in the right panel of Figure 6) exhibit markedly different behavior. Their sizes remain significantly larger than both the global sizes of QGs (black curves) and the sizes of QG bulges (orange/maroon curves). Although QG bulges grow rapidly with time, the growth rate of their disks — unlike in SFGs — is comparable to that of their global sizes. These highlight key structural differences in how SFGs and QGs grow over cosmic time.

While we compare the global size evolution from Paper I with the component sizes derived in this work, we caution the reader that these samples are not strictly identical. Specifically, the galaxy populations contributing to the global curves (black) and the component curves (colored) in Figure 6 differ in their selection criteria. Furthermore, the sub-samples used for bulge and disk evolution differ from one another (see Section 4 and Table 1). Consequently, the global size evolution is not a simple superposition of the bulge and disk evolutionary tracks, particularly for SFGs where selection effects significantly impact the composition of the bulge and disk samples.

As noted in Section 5, the pace of structural evolution exhibits a clear wavelength dependence. For both SFGs and QGs, bulge components show more rapid expansion in the rest-UV than in the optical; specifically, between $z = 0.85$ and $z = 0.1$ (~ 6 Gyr), QG (SFG) bulges grew by 120% (98%) in the UV compared to a more modest 76% (44%) in the optical. Conversely, disk growth in QGs is more pronounced in the optical (68%) than in the UV (29%). Notably, SFG disks show no significant difference between their UV (32%) and optical (40%) growth rates. This suggests that while bulge and quiescent disk evolution are sensitive to wavelength-dependent stellar population gradients, SFG disks likely grow through more spatially uniform, continuous star formation across the sampled wavelengths.

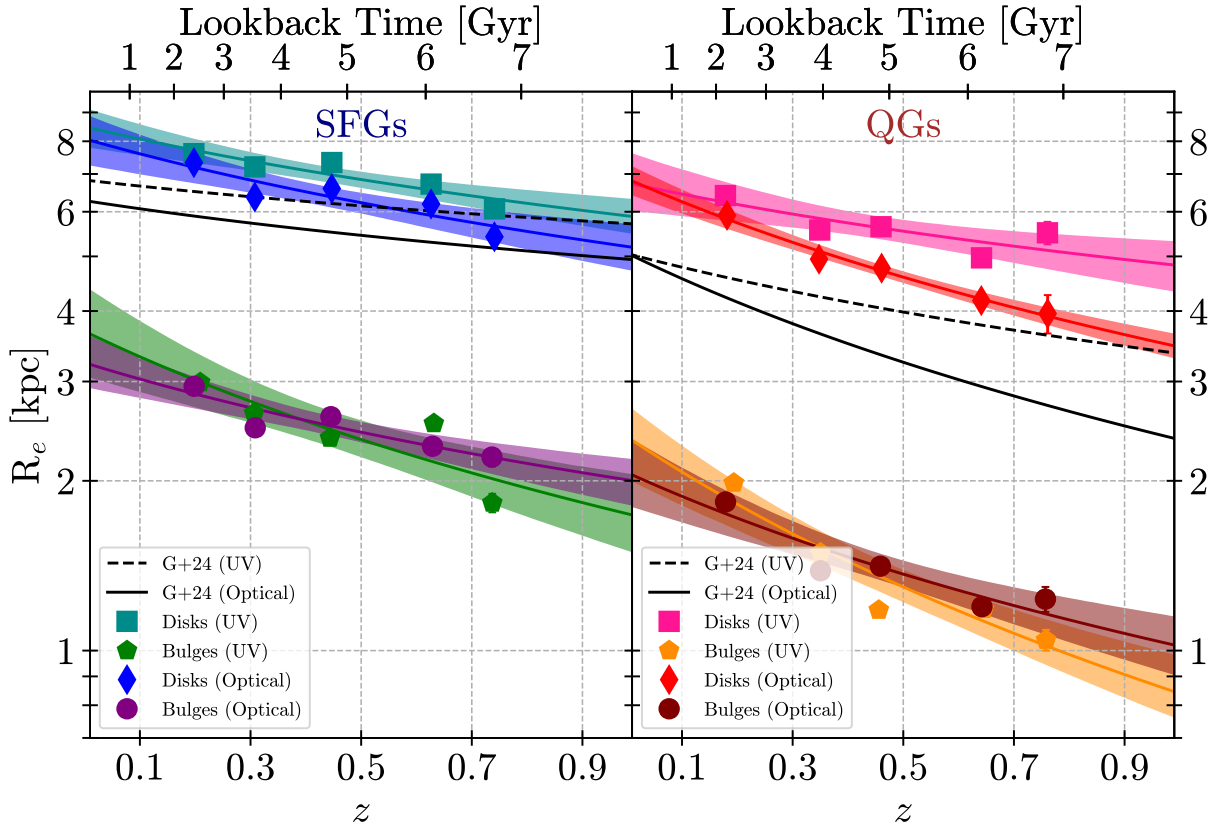


Figure 6. Left: Redshift evolution of characteristic sizes of SFG bulges and disks in the rest-frame UV (3000Å) and optical (5000Å) wavelengths. The global size evolution curves from A. George et al. (2024, G+24) based on single Sérsic profile fitting are shown in black. Right: Redshift evolution of characteristic sizes of QG bulges and disks in two rest-frame wavelengths.

899 Finally, for both SFGs and QGs, bulges exhibit a
 900 faster rate of size growth compared to disks. This trend
 901 is consistent with recent findings by K. V. Nedkova et al.
 902 (2024), who also reported more rapid growth in the char-
 903 acteristic sizes of bulges relative to disks based on the
 904 Hubble Frontier Fields and CANDELS data. In our
 905 analysis, we further find that this difference in growth
 906 rates between bulges and disks is wavelength-dependent.
 907 Specifically, the disparity is more pronounced at rest-
 908 frame 3000Å than at 5000Å: Growth curve is 105%
 909 (208%) steeper for SFG (QG) bulges than their disks in
 910 rest-UV, but the difference is negligible — 9% (3%) — in
 911 rest-optical. Thus, the UV light more clearly reveals the
 912 differential size evolution between galaxy components
 913 than optical light.

914 Synthesizing these trends, we find that the structural
 915 growth is fundamentally coupled to both the specific
 916 galaxy component and the observed wavelength. Figure
 917 7 presents a schematic of an average two-component
 918 SFG (left) and QG (right) with a fiducial stellar mass
 919 $\log M \sim 10.7$ at $z \sim 0.5$. The disks of both SFGs
 920 and QGs appear more extended in rest-frame UV light
 921 compared to the optical light. Meanwhile, the bulges are

922 slightly larger, if not the same, in the optical than in
 923 the UV. Moreover, the difference between the UV sizes
 924 of bulges and disks is more pronounced in QGs than in
 925 SFGs. These wavelength-dependent differences in the
 926 sizes along with the differences in their redshift evolu-
 927 tion (Figure 6) highlight the distinct physical processes
 928 shaping bulges and disks, and underscore the role of stel-
 929 lar population gradients and structural growth in galaxy
 930 morphological evolution.

931 7. DRIVERS OF BULGE AND DISK SIZE 932 EVOLUTION

933 Our bulge+disk decomposition of the CLAUDS+HSC
 934 galaxies across two rest-frame wavelengths confirms that
 935 bulges are significantly more compact than their com-
 936 panion disks and, crucially, reveals that these structural
 937 components follow distinct evolutionary pathways. Be-
 938 yond the wavelength-dependence of the sizes themselves,
 939 the pace of structural evolution also varies with wave-
 940 length and galaxy type (Section 6). These nuances sug-
 941 gest that while certain global trends persist for all bulges
 942 and disks, their specific growth trajectories are heav-
 943 ily influenced by the star-formation activity of the host

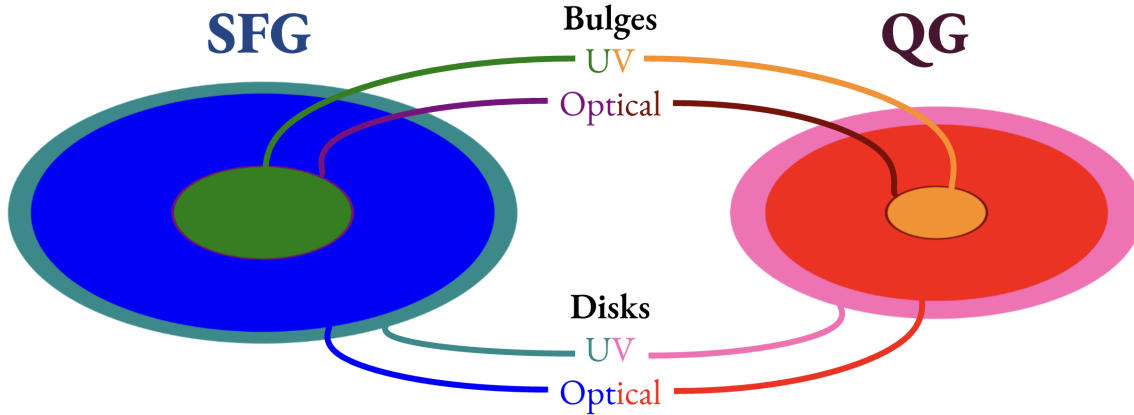


Figure 7. A simple model of an average SFG (left) and QG (right) with a fiducial stellar mass $\log M \sim 10.7$ at $z = 0.5$. The sizes of the disks and bulges are shown to scale relative to each other. Disks appear more extended in the rest-UV (cyan/pink) than in the rest-optical (blue/red), whereas bulges are more compact in the UV (green/orange) compared to the optical (purple/maroon). Additionally, both bulges and disks are smaller in QGs than in SFGs. For comparison, a constant axis ratio (q) is adopted for both galaxy types.

944 galaxy. Such divergence warrants a separate examination of the physical mechanisms driving the structural
 945 evolution of the star-forming and quiescent populations.
 946

947 7.1. Star-forming Galaxies

948 The simple schematic picture of an SFG with the fidu-
 949 cial mass observed in the multi-wavelength analysis us-
 950 ing bulge+disk decomposition (Figure 7) and the trends
 951 observed in the evolution of structural parameters (the
 952 left panel of Figure 6) support the evolutionary scenario
 953 proposed in our Paper I. There, we concluded that the
 954 observed overall size evolution of the SFGs (from single-
 955 component Sérsic fits at two rest-frame wavelengths)
 956 must be shaped by the growth of bulges within them.
 957 The fast size growth of bulges that we observe in this
 958 work provide additional arguments for our earlier inter-
 959 pretation.

960 Bulges are centrally concentrated stellar distributions
 961 in galaxies, that become prominent in galaxies with de-
 962 creasing redshift (M. Martig et al. 2009; S. Sachdeva
 963 2013; P. Lang et al. 2014; S. Sachdeva et al. 2017; S. Tac-
 964 chella et al. 2019; A. George et al. 2024). There are sev-
 965 eral physical mechanisms that can drive their formation
 966 and growth. Classical bulges, which exhibit properties
 967 similar to early-type galaxies, can initially form through
 968 the dissipative collapse of protogalactic gas clouds (O. J.
 969 Eggen et al. 1962) or the coalescence of giant clumps in
 970 primordial disks (e.g., F. Bournaud et al. 2007; N. M.
 971 Förster Schreiber & S. Wuyts 2020). These structures
 972 can subsequently grow through secular disk instabilities
 973 (G. De Lucia et al. 2011) or hierarchical galaxy mergers
 974 (P. F. Hopkins et al. 2010a; A. Brooks & C. Christensen
 975 2016).

976 Galaxy mergers are thought to play a crucial role in
 977 the growth of bulges in galaxies. Though major mergers
 978 can form central bulges and form/reform disks around
 979 them from the surviving material (J. E. Barnes & L.
 980 Hernquist 1992, 1996; V. Springel & L. Hernquist 2005;
 981 T. Naab et al. 2006; P. F. Hopkins et al. 2009), such
 982 mergers may not have played a crucial role in bulge for-
 983 mation/growth in the last 7 billion years due to their low
 984 frequency (C. J. Conselice et al. 2003, 2022; C. López-
 985 Sanjuan et al. 2009, 2010; J. M. Lotz et al. 2011; S.
 986 Kaviraj et al. 2014; N. Thibert et al. 2021; Q. Duan
 987 et al. 2025). However, during this epoch, minor mergers
 988 are frequent (J. M. Lotz et al. 2011) and can signifi-
 989 cantly facilitate bulge growth through the redistribu-
 990 tion of mass and angular momentum. In these events,
 991 the satellite galaxy is disrupted by tidal forces, with its
 992 stellar and gaseous content sinking toward the potential
 993 well of the host via dynamical friction (M. Steinmetz
 994 & J. F. Navarro 2002; L. Oser et al. 2010). Simultane-
 995 ously, the gravitational interaction kinematically heats
 996 the host disk, scattering stars from ordered circular or-
 997 bits into the random, pressure-supported orbits charac-
 998 teristic of a bulge component (T. J. Cox et al. 2008;
 999 P. F. Hopkins et al. 2010b). This combination of di-
 1000 rect satellite accretion and the dynamical restructuring
 1001 of the host disk provides a robust mechanism for the
 1002 continuous assembly of bulges since $z \sim 0.85$.

1003 Bulges can also grow through internal (secular) pro-
 1004 cesses, which differ from merger-driven growth in both
 1005 their timescale and dynamical nature. Unlike the violent
 1006 kinematic heating caused by mergers, secular evolution
 1007 is a gradual process where internal disk instabilities —

1008 such as bars or spiral arms — redistribute angular mo- 1060
 1009 mentum within the disk plane (E. Athanassoula 2005; 1061
 1010 J. A. Sellwood 2014). These structures funnel gas and 1062
 1011 stars toward the center without destroying the disk’s 1063
 1012 underlying rotational support, allowing mass to settle 1064
 1013 into a more stable, centrally concentrated configuration 1065
 1014 (J. Kormendy & D. B. Fisher 2008; L. Zhang et al. 1066
 1015 2025). The resulting pseudobulges retain disk-like kine- 1067
 1016 matics and nearly exponential light profiles, distinguish- 1068
 1017 ing them from the pressure-supported classical bulges 1069
 1018 formed through external interactions (J. Kormendy & 1070
 1019 R. C. Kennicutt 2004; D. B. Fisher & N. Drory 2011; J. 1071
 1020 Hu et al. 2024).

1021 Regardless of the underlying growth mechanism, the 1073
 1022 emergence of a bulge fundamentally modulates both the 1074
 1023 structural morphology and the subsequent evolutionary 1075
 1024 trajectory of the host galaxy. Due to their high cen- 1076
 1025 tral stellar mass densities and compact profiles, bulges 1077
 1026 shift the global light distribution toward more concen- 1078
 1027 trated states, effectively reducing the measured global 1079
 1028 effective radius. Beyond these geometric alterations, 1080
 1029 the development of a substantial central mass compo- 1081
 1030 nent can kinematically stabilize the surrounding gas disk 1082
 1031 against gravitational fragmentation (M. Martig et al. 1083
 1032 2009). This stabilization regulates the efficiency of star 1084
 1033 formation by increasing the marginal stability of the 1085
 1034 disk, suggesting that bulge assembly is a primary driver 1086
 1035 in steering a galaxy’s progression along the star-forming 1087
 1036 sequence toward eventual quiescence.

1037 Building on our earlier results from Paper I, we pro- 1089
 1038 posed two main scenarios for bulge growth in SFGs. In 1090
 1039 Scenario I, both the bulge and the disk grow concu- 1091
 1040 rrently as the SFG evolves. In Scenario II, the bulge 1092
 1041 alone grows, and this growth pushes the active star- 1093
 1042 forming regions farther out in the disk. We can probe 1094
 1043 these two scenarios in more detail with the results from 1095
 1044 our multi-wavelength bulge+disk decomposition.

1045 Our results show that SFG disks are consistently more 1097
 1046 extended in the rest-UV than in the optical, with both 1098
 1047 radii increasing over cosmic time (Figure 6). While the 1099
 1048 larger rest-UV sizes are indicative of inside-out growth 1100
 1049 — where younger stellar populations are preferentially 1101
 1050 located in the disk outskirts — this observation is con- 1102
 1051 sistent with both scenarios. The primary differentia- 1103
 1052 tor is the evolution of the bulge-to-total (B/T) ratio, 1104
 1053 which remains relatively stable across our redshift range. 1105
 1054 Although the B/T ratio is inherently more difficult to 1106
 1055 constrain than the effective radius due to the degener- 1107
 1056 acies in light-profile modeling (Section 4), the lack of 1108
 1057 significant B/T evolution suggests that the disk is also 1109
 1058 growing concurrently. This indicates that while bulges 1110
 1059 expand in physical size more rapidly than disks, their

1060 relative light contribution remains relatively steady, pro-
 1061 viding the strongest evidence for the concurrent growth
 1062 described in Scenario I.

1063 Furthermore, the sizes of SFG disks grow faster than
 1064 the total (single Sérsic) sizes of SFGs themselves in
 1065 both rest-frame wavelengths (the left panel of Figure 6).
 1066 The total sizes of SFGs and the sizes of their disks are
 1067 very similar in the highest redshift bin, but the total
 1068 sizes of SFGs deviate significantly from the disk size
 1069 with decreasing redshift. The observed deceleration in
 1070 the growth of global sizes of the SFGs at $z < 1$ from
 1071 our pilot study must be driven by the size growth of
 1072 bulges. In fact, the growth rate of disks in rest-optical
 1073 ($\beta = -0.64 \pm 0.24$) we find here is consistent with the
 1074 fast growth of SFGs reported at higher redshifts by R. J.
 1075 Williams et al. (2010, $\beta = -0.77$ at $0.5 < z < 2$) and A.
 1076 van der Wel et al. (2014, $\beta = -0.75$ at $0 < z < 3$), sup-
 1077 porting our hypothesis. This demonstrates that while
 1078 the physical growth of the disk component remains rapid
 1079 and consistent with high-redshift trends, the apparent
 1080 deceleration of the global size growth is a measurement
 1081 effect driven by the increasing influence of the compact
 1082 bulge on the single-Sérsic fit. Specifically, as the bulge
 1083 becomes a more structurally distinct feature at lower
 1084 redshifts, it increases the global Sérsic index (n) of the
 1085 galaxy; this sharpens the central light profile and an-
 1086 chors the global R_e to smaller values, effectively masking
 1087 the true extent of the disk’s expansion.

1088 The size growth of the SFG disks is slightly faster
 1089 in the longer wavelength ($\beta = -0.64 \pm 0.24$) than in
 1090 the shorter wavelength ($\beta = -0.53 \pm 0.19$). Although
 1091 within uncertainties, this small difference in the growth
 1092 rate could indicate that the star-forming activity in disks
 1093 may be slowing down as the SFGs evolve. In such a
 1094 scenario, we would expect a slower size growth in rest-
 1095 UV than in rest-optical, as observations suggest.

1096 This gradual decline in the star-formation activity of
 1097 SFGs over cosmic time can be linked to the growth of
 1098 their internal bulges. By kinematically stabilizing the
 1099 surrounding gas disk, the galactic bulges can suppress
 1100 the star formation efficiency, a process known as mor-
 1101 phological quenching (M. Martig et al. 2009; A. Sain-
 1102 tonge et al. 2012; S. Sachdeva et al. 2017; A. Hashem-
 1103 izadeh et al. 2022). This interpretation is consistent
 1104 with the observed decline in the specific star formation
 1105 rate (sSFR) with increasing stellar mass surface density
 1106 within the central 1 kpc radius (Σ_1 in units of M_\odot/kpc^2 ;
 1107 Paper I). As bulge growth drives Σ_1 towards a thresh-
 1108 old value ($\log \Sigma_1 \sim 9$), the sSFR begins to decrease via
 1109 inside-out suppression of star formation (K. E. Whitaker
 1110 et al. 2017; A. George et al. 2024). Further bulge growth
 1111 can eventually quench a galaxy; galaxies with very high

central mass density ($\log \Sigma_1 \gtrsim 9.7$) are generally quiescent (K. E. Whitaker et al. 2017; K. A. Suess et al. 2021; V. Estrada-Carpenter et al. 2023; A. George et al. 2024). The physical evidence for this process is also reflected in our finding that B/T values are higher in the optical than in the UV (Figure 3).

However, we observe that bulges grow faster in UV, suggesting an ongoing addition of UV-bright material. Since we model bulges with a free Sérsic index, the fits may include other non-disk components, such as stellar halos. Hence, our bulge measurements may include material accreted through minor mergers, which tend to settle in the outer regions, contributing to halo growth (D. J. Williams et al. 2025). These accreted *ex-situ* stars can be younger and/or lower in metallicity, making them more prominent in the UV and thus boosting UV-measured sizes more than optical ones (I. Saviane et al. 2008; A. P. Cooper et al. 2010; M. Hirschmann et al. 2015).

The gradual increase in the central concentration of bulge components is reflected in the distribution of bulge Sérsic indices (n_B 's) in SFGs and their evolution with redshift (upper panel of Figure 8). We find that the Sérsic indices of SFG bulges in the optical light (purple) increase over cosmic time, such that by $z < 0.4$, their profiles in the rest-frame optical approach a de Vaucouleurs profile ($n = 4$), consistent with trends reported by P. Dimauro et al. (2019). This suggests that the structural properties of present-day SFG bulges in the rest-optical resemble those of QG bulges, indicating a gradual morphological transformation of bulges toward early-type characteristics. However, we do not find any significant evolution occurring in the UV light (green), where $n_B \sim 1$. This indicates that the distribution of star-forming regions and/or metal-poor stars within SFG bulges remains exponential over the redshift interval we probe.

In summary, we find strong evidence for the growth of centrally concentrated bulges that play a vital role in the evolution of SFGs since $z < 1$. Although this growth of bulges could be driven by various processes (e.g., *in-situ* star formation, disk instability, mergers, etc.), the process rapidly builds up materials in galaxy center and promotes inside-out quenching.

7.2. Quiescent Galaxies

Although general trends in the size evolution of QGs (such as relative sizes of components, their size growth, and wavelength dependence) are similar to those of SFGs, there are some notable differences. The sizes of the disks and bulges of QGs are smaller than those of SFGs (Sections 5 and 6). In addition, the global (single-

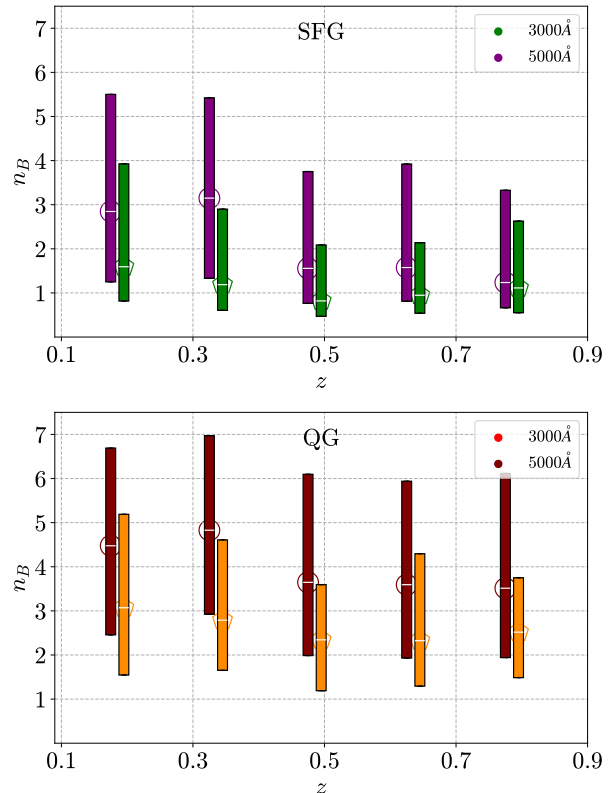


Figure 8. Top: Evolution of Sérsic indices of SFG bulges with stellar mass $10.5 < \log M < 11$ in the rest-UV and in the rest-optical. The colored boxes extend from the first quartile to the third quartile of the Sérsic indices in different redshift bins, with a white horizontal line at the median value. Bottom: Evolution of Sérsic indices of QG bulges with stellar mass $10.5 < \log M < 11$ in the rest-UV and in the rest-optical. Rest of the details are same as in the top panel. We also note that SFG and QG samples are not mass complete at $z > 0.7$ and at $z > 0.55$ respectively.

Sérsic) sizes of QGs are significantly different from the sizes of their individual components. The pace of the size evolution for QG disks is comparable to that of SFGs in the rest-UV. However, the QG disks (in optical) and bulges (in UV and optical) grow at a faster rate than their SFG counterparts.

A significant fraction of QGs retain a disk or a component that can be modeled with an exponential light profile. This is in agreement with previous observational works showing that disk structures exist in many if not all QGs (e.g., P. Dimauro et al. 2019; Q. Cui et al. 2024; K. V. Nedkova et al. 2024; F. Rigamonti et al. 2024). Survival of disks after quenching is also predicted in cosmological simulations. For example, the analysis of galaxy kinematic properties from the EAGLE cosmological simulation shows that when star formation in massive galaxies ($\log M > 10.5$) declines, they trans-

1180 form into spheroids with a disk component (B. Clauwens
1181 et al. 2018). These transformations are driven mainly
1182 by the accretion of material in the outer halo, allowing
1183 disks to survive — albeit in a faded state. By trac-
1184 ing the evolution of the SMR for QG disks in two rest-
1185 frame wavelengths, our analysis provides a unique ob-
1186 servational probe of how these quenched disks evolve.

1187 Compared to bulges, disks in QGs are not compact
1188 (the right panel of Figure 6). Their sizes are compara-
1189 ble to those of the disk components of the SFG popu-
1190 lation. This similarity in disk sizes between SFGs and
1191 QGs could indicate that disk formation in galaxies oc-
1192 curs while they are still active. In Paper I, we find that
1193 20% of the size growth of QGs comes from the progen-
1194 itor bias driven by newly quenched galaxies (newcom-
1195 ers), a result recently corroborated across both rest-UV
1196 and optical wavelengths by M. Pizzardo et al. (2026).
1197 Some of the progenitor bias effects should be reflected
1198 in the disk profiles, as these newcomer galaxies tend
1199 to maintain more prominent and extended disk struc-
1200 tures than those quenched at earlier epochs (M. Piz-
1201 zardo et al. 2026). Additionally, we observe that qui-
1202 escent disks grow faster in the optical wavelength than
1203 in the UV. Thus, if QGs have a disk component, these
1204 disks are more extended in UV than in the visible light
1205 at higher redshifts, but there is no significant size dif-
1206 ference between wavelengths at low redshift due to the
1207 slower growth in the UV.

1208 We should also note that, although the disk fades
1209 when a galaxy quenches (e.g., D. Christlein & A. I.
1210 Zabudoff 2004; C. M. Carollo et al. 2016; J. Matharu
1211 et al. 2020; V. Estrada-Carpenter et al. 2023), this fad-
1212 ing happens over several hundreds of millions of years.
1213 In fact, our rest-frame 3000 Å size measurements are
1214 highly sensitive to the age of the stellar population. At
1215 this wavelength, the light is dominated by stars formed
1216 over the last several hundred million years; specifically,
1217 the average flux-weighted age of a simple stellar popu-
1218 lation with solar metallicity is ~ 300 Myr (G. Bruzual
1219 & S. Charlot 2003; D. C. Martin et al. 2005). Thus, the
1220 contribution from the newcomers should be more promi-
1221 nent in the rest-UV than in the longer wavelength, as
1222 evidenced by the similar UV size evolution slopes (β)
1223 observed between SFG and QG disks.

1224 However, QG disks exhibit a significantly higher β in
1225 the optical than their SFG counterparts. This accel-
1226 erated optical growth indicates that QG disks are not
1227 merely fading SFG disks; instead, they are subject to
1228 cumulative structural assembly via minor mergers that
1229 deposit mass into the outskirts. While these mergers
1230 can deliver younger/bluer stars — a process that drives
1231 the rapid UV growth seen in QG bulges — the overall

1232 disk size evolution in the UV is ultimately limited by the
1233 rapid fading of these young populations once they are
1234 incorporated. Consequently, while the UV light traces
1235 the ‘memory’ of recent star formation in newcomer disks,
1236 the optical growth provides a more permanent record of
1237 the structural assembly and stellar mass redistribution
1238 occurring over longer timescales.

1239 The bulges in QGs are more compact than the SFG
1240 bulges (the right panel of Figure 6) and have higher
1241 values of n_B (the bottom panel of Figure 8) in both
1242 rest-frame wavelengths, indicating that QG bulges have
1243 significantly more central concentration. The lack of
1244 significant evolution in n_B for QGs with redshift sug-
1245 gests that quenching mechanisms at $z < 1$ are closely
1246 coupled with morphological transformation, typically re-
1247 sulting in de Vaucouleurs ($n = 4$) bulges in rest-optical
1248 (maroon bars in the bottom panel of Figure 8). While
1249 newcomers account for approximately 20% of the popu-
1250 lation’s size growth (Paper I), their inclusion does not
1251 shift the median n_B significantly, implying that these
1252 galaxies already have significant centrally concentrated
1253 profiles when they join the quiescent population.

1254 As galaxies quench and move through the green valley,
1255 they undergo significant morphological transformations.
1256 In the optical regime, these transitions are character-
1257 ized by a decrease in R_e and an increase in n , lead-
1258 ing to higher stellar mass surface densities without an
1259 appreciable increase in total stellar mass (L. Quilley
& V. de Lapparent 2022; V. Estrada-Carpenter et al.
1260 2023). However, structural evolution may continue be-
1261 yond the initial quenching phase. Post-quenching, QGs
1262 can become even more centrally concentrated over time,
1263 mainly due to the continued accumulation of stellar mass
1264 through minor mergers and accretion (e.g., M. Hilz et al.
1265 2013; H. J. Zahid et al. 2019). Even though QG bulges
1266 typically do not exhibit de Vaucouleurs profiles in the
1267 shorter wavelength (orange bars), their n_B values are
1268 consistently higher than those of SFG bulges in the
1269 same wavelength, supporting the notion of morphologi-
1270 cal transformations during and after quenching.

1271 Moreover, the bulges of the quenched galaxies are
1272 slightly more compact in UV than in the optical at
1273 $z > 0.4$ (the right panel of Figure 6). Although these
1274 compact UV bulges could indicate residual star forma-
1275 tion in the cores (C. R. Wagner et al. 2015; A. W. S. Man
1276 et al. 2016; C. D’Eugenio et al. 2020), galaxy centers are
1277 typically older than the galaxy outskirts (K. A. Suess
1278 et al. 2019, 2020). Fast processes that form bulges at
1279 high z ’s (e.g., wet compaction) could produce some low-
1280 metallicity stars in the centers of galaxy bulges (R. L.
1281 Griffith et al. 2011; A. R. Casey & K. C. Schlaufman
1282 2015; S. Lapiner et al. 2023), even though galaxies in
1283

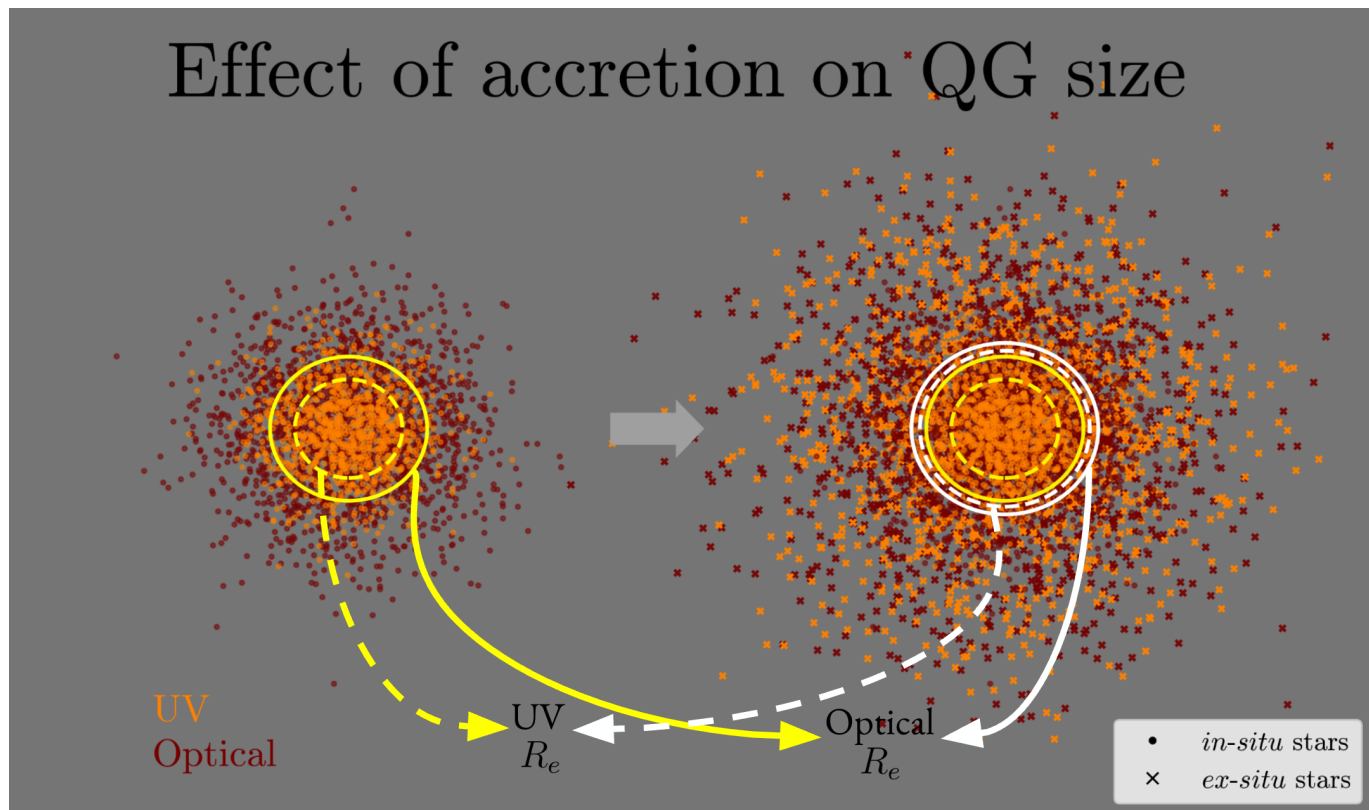


Figure 9. A schematic diagram showing the effect of minor mergers/accretion on the size of bulge components of a QG in the rest-frame UV (orange) and optical (maroon) wavelengths. The effective radii before and after adding materials to the galaxy are shown in yellow and white, respectively. The materials added through minor mergers/accretion occupy the outer halo of a galaxy thereby increasing the measured effective radius of its bulge component. Unlike the *in-situ* stellar population (dots) of a QG, the newly added *ex-situ* population (crosses) have a significant fraction of UV bright stars (young massive stars or low metallicity stars). Thus, the merger/accretion-driven growth is faster in rest-UV than in the rest-optical.

1284 general show negative metallicity gradients. Studies 1306
 1285 indicate that some of the galaxy bulges are relics of 1307
 1286 blue/red nugget galaxies from high redshifts ($z > 1$; I. 1308
 1287 Damjanov et al. 2009, 2011; A. Ferré-Mateu et al. 2017; 1309
 1288 L. Costantin et al. 2021, 2022; C. Spiniello et al. 2021). 1310
 1289 During the post-blue nugget phase, these galaxies tend 1311
 1290 to have a flat metallicity gradient in the innermost 1 kpc 1312
 1291 region due to fast depletion of gas (S. Lapiner et al. 1313
 1292 2023), and they form a major part of the bulges of more 1314
 1293 evolved galaxies. The presence of similar low-mass stars 1315
 1294 in the bulge of the Milky Way has been well explored 1316
 1295 (L. M. Howes et al. 2015; M. Schultheis et al. 2015; M. 1317
 1296 Ness & K. Freeman 2016; F. Fragkoudi et al. 2017; H. 1318
 1297 Reggiani et al. 2020). Because the UV-optical color of 1319
 1298 low metallicity stars is bluer than that of high metallic- 1320
 1299 ity stars, their presence should make galaxy bulges more 1321
 1300 compact in the rest-UV than in the rest-optical. 1322

1301 Additionally, the Sérsic indices of QG bulges are lower 1323
 1302 in the UV than in the optical (lower panel of Figure 8). 1324
 1303 This indicates that while the UV-emitting stars occupy 1325
 1304 a compact region, their spatial distribution is less centrally 1326
 1305 concentrated than that of the bulk stellar popu- 1327

1306 lation. This result is consistent with the recent findings
 1307 of M. Pizzardo et al. (2026), who utilized single Sérsic
 1308 profile fits to show that quiescent galaxies exhibit lower
 1309 n values in the rest-UV compared to the optical.

1310 This less centrally peaked distribution in UV could
 1311 arise for multiple reasons. The UV profile could reflect
 1312 the low n distribution of the star-forming regions prior to
 1313 quenching. They could also be driven by the distribution
 1314 of globular clusters with UV bright or low metallicity
 1315 stars (e.g., S. Kaviraj et al. 2007; D. A. Forbes et al.
 1316 2022).

1317 Moreover, shallow UV profiles could also be the result
 1318 of the distribution of young/low-metallicity stars in the
 1319 outskirts added through minor mergers and accretion
 1320 (halo growth; D. J. Williams et al. 2025). While minor
 1321 mergers are understood to increase the global Sérsic in-
 1322 dex in terms of stellar mass by building extended wings
 1323 in the mass distribution (e.g., P. F. Hopkins et al. 2010a;
 1324 H. J. Zahid et al. 2019), they may have the opposite ef-
 1325 fect on the UV light profile. If the accreted material is
 1326 more UV-luminous than the central population — due
 1327 to downsizing in time (E. Neistein et al. 2006) or low

metallicity — its presence in the halo would ‘flatten’ the light distribution. Because our bulge model captures all non-disk light (Section 7.1), including this accreted halo, the combination of a central bulge and a UV-bright outer halo would yield a lower n in the UV than in the optical, where the older, centrally-concentrated mass dominates.

Furthermore, the sizes of QG bulges grow rapidly with cosmic time. This growth rate is higher in UV light than in visible light; although QG bulges are slightly smaller in UV than in the optical at the high redshift limit of our sample, the UV sizes become slightly larger than optical sizes at $z < 0.4$ (the right panel of Figure 6). The wavelength dependence of the size increase for bulges supports the minor merger-/accretion-driven growth of QGs because the accreted material has a higher fraction of UV bright stars (young or low-metallicity). Since the QGs are in general UV faint, the addition of accreted material to the galaxy halos increases the effective radii of bulges in the UV more than in the optical.

We illustrate the impact of *ex-situ* accreted material in QGs on the effective radius measured at the rest-UV (orange) and rest-optical (maroon) wavelengths using a simplified toy model in Figure 9. The image on the left is a representative diagram of QG bulges in these two wavelengths. There are more maroon points than orange ones, indicating that the bulge is brighter in the longer wavelengths¹⁵. The solid and dashed curves represent the measured effective radii in the UV and optical wavelengths, respectively, with the bulge appearing more compact in the UV than in the optical.

When this galaxy undergoes a series of minor mergers or accretion events, material is added to the outer regions. The accreted population in this toy model consists of an equal proportion of both orange and maroon data points, as illustrated in the galaxy model on the right. This is because the accreted material comes from low-mass systems could contain higher fractions of younger/low-metallicity stars than massive quiescent systems (I. Saviane et al. 2008; A. P. Cooper et al. 2010; M. Hirschmann et al. 2015). The addition of stellar material changes the effective radius of orange points (dashed circles) more than that of the maroon points (solid circles), *i.e.*, the measured size grows more in UV light than in the optical light. Therefore, the growth of bulge/spheroidal components driven by minor mergers and accretion is faster in the shorter wavelength than in the longer wavelength.

¹⁵ Note that a significant fraction of maroon points is behind the distribution of orange points

In addition, part of the observed growth in the QG bulge sizes may result from progenitor bias introduced by the newcomers. At any given redshift, SFG bulges tend to be larger than QG bulges, so the ongoing addition of these newly quenched galaxies to the QG population should increase the average quiescent bulge size. To summarize, as in the case of the total sizes of QGs, the size evolution of their structural components is also driven by mergers/accretion along with the progenitor bias from the newcomers.

8. LIMITATIONS

We note that the limitations we discuss in our Paper I are also valid for this work. Classification of galaxies into star-forming and quiescent depends on the methodologies used (color-color diagrams, specific star-formation rates, etc.) and could differ from other works in the literature. Using our data set, we cannot distinguish between contributions from young and low-metallicity stars in the UV light profiles (due to the age-metallicity degeneracy; B. Gustafsson 1989; D. Streich et al. 2014). The presence of dust can also affect measured sizes, particularly in the rest-UV band, where the dust obscuration is higher than that of the visible band. Our dataset cannot trace the radial distribution of dust in our targets; spatially resolved data at longer wavelengths (far-infrared to radio) is required for dust mapping.

In addition, there are limitations to the fitting methodology adopted in this work. While we use an exponential profile for the disk and a free Sérsic profile for the bulge, many massive galaxies may require a third “halo” component to fully capture their outer stellar envelopes. Although our deep ground-based data theoretically reaches the surface brightness levels required to detect such halos, robustly decomposing three separate components at these redshifts remains challenging. The primary obstacle is model degeneracy; the spatial resolution of ground-based imaging often cannot distinguish between a concentrated bulge and an extended halo component, particularly when they overlap in the central regions.

To manage these degeneracies, some studies fix the bulge to a de Vaucouleurs profile ($n = 4$). While this reduces the uncertainty in size and flux measurements, it imposes a structural bias by forcing all bulges into a centrally concentrated shape regardless of their intrinsic profile. If a galaxy possesses a lower-Sérsic index bulge, fixing $n = 4$ results in significant systematic residuals and can lead to inaccurate measurements of the bulge effective radius and B/T values. By allowing n to remain free, we prioritize capturing the intrinsic variety

of bulge shapes, even if it introduces higher statistical uncertainty.

Moreover, our selection of disks and bulges, based on B/T cuts (Section 4), excludes components that contribute less than 20% to the galaxy’s total surface brightness profile. By implementing this threshold, we ensure that our analysis is restricted to galaxies with two distinct and well-defined structural components. This prevents the inclusion of marginal ‘bulges’ or ‘disks’ — which may simply represent noise or minor residuals in a single-component system — thereby allowing for a more robust comparison of the physical evolution of these two distinct populations.

Furthermore, light contamination from nearby faint background objects can affect component size measurements, particularly for intrinsically faint bulges/disks (Section 3). Our simulations show that such contamination tends to increase measured sizes. While we mitigate this by excluding very faint components through B/T cuts (Section 5) and introduce a correction to the measured sizes using simulations (Appendix A), the effect is not entirely removed. For instance, the large optical sizes of QG disks ($\log R_e \sim 1.2$ at $\log M \sim 11$) in the bottom panel of Figure 4 are partly due to this effect. However, we do not expect this to bias our scientific conclusions, as we account for such uncertainties through the error estimates described in Appendix A.

Finally, the SMRs of bulges and disks strongly depend on the choice of the stellar mass (J. Méndez-Abreu et al. 2021). Several works in the literature adopt the stellar mass of the components to study their scaling relation (e.g., P. Dimauro et al. 2019; A. S. G. Robotham et al. 2022b; K. V. Nedkova et al. 2024). However, estimating the stellar masses of individual components is far more challenging than measuring the total stellar mass. SED-based mass estimates depend sensitively on assumptions about the stellar population models, IMF, and dust attenuation laws (C. Pacifici et al. 2023; K. V. Nedkova et al. 2024), and separating the bulge and disk stellar populations is not straightforward. In addition, the IRAC data used to constrain the mass estimates have a relatively poor PSF, which further limits our ability to obtain reliable component masses. Therefore, we adopt the total stellar mass while analyzing the SMR for bulges and disks. As we discuss in Section 5, in this work, we probe the evolution of SMR-derived characteristic sizes of bulges and disks in galaxies with similar stellar mass.

9. CONCLUSIONS

Utilizing deep imaging from the CFHT CLAUDS and Subaru HSC-SSP surveys, we have characterized the structural evolution of a mass-complete sample compris-

ing 44,086 SFGs and 34,108 QGs across the redshift range $0.1 < z < 0.85$. By performing bulge+disk decompositions with SOURCEEXTRACTOR++ in both the rest-frame UV (3000Å) and optical (5000Å), we examine the size–mass relation (SMR) for the bulge and disk components separately. From these relations, we determine the characteristic size of each component for galaxies at a fiducial total stellar mass of $5 \times 10^{10} M_\odot$, allowing us to track how the physical scales of these structures evolve with cosmic time. This comparison probes different aspects of bulge and disk growth for the two galaxy populations (Figure 6):

1. *Relative component scales* - The characteristic sizes of disks are consistently larger than those of bulges for both SFGs and QGs across both wavelengths.
2. *Wavelength-dependent sizes* - For a given galaxy population, disk sizes are more extended in the UV than in the optical, whereas bulge sizes remain comparable between the two wavelengths.
3. *Differential growth rates* - The pace of size evolution varies significantly between populations and wavelengths. Bulges exhibit faster size evolution in the rest-UV than in the rest-optical. In contrast, QG disks grow significantly faster in the longer wavelength than in the shorter wavelength, whereas this wavelength-dependent difference is much less pronounced for SFG disks.

Thus, our results show that if a galaxy has both a bulge and a disk, its disk would be larger than its bulge and would appear more extended in UV than in the optical (Figure 7). In contrast, the bulges are slightly smaller in UV light at higher z , but due to their fast growth, the UV size catches up with the optical size by $z = 0.4$.

These results confirm our previous finding based on single-Sérsic fitting (Paper I) that bulge growth is a major driving mechanism in SFGs. The bulges are growing at a rapid rate, and their growth can affect the measured sizes of the SFGs using a single Sérsic profile. The emergence/growth of bulges could significantly slow the size growth rate of SFGs, although the individual structural components (bulge and disk) grow faster than their combined size.

We also interpret the rapid size growth of QG bulges (more in the UV band than in the optical band) as driven by minor mergers and accretion. This merger-driven growth affects more in the shorter wavelength than in the longer wavelength (Figure 9).

1526 However, we cannot currently disentangle the effects
 1527 of the age-metallicity degeneracy and impact of dust
 1528 with our dataset. We need spatially resolved spec-
 1529 troscopy and dust maps to address these issues. Ongo-
 1530 ing and upcoming large-scale surveys using cutting-edge
 1531 integral field spectrographs/units like Very Large Tele-
 1532 scope’s MUSE (T. Urrutia et al. 2019; C. Foster et al.
 1533 2021; E. Emsellem et al. 2022) and adaptive optics (AO)
 1534 assisted ERIS (R. Davies et al. 2018; J. Hayoz et al.
 1535 2025), James Webb Space Telescope’s NIRSpec IFU (T.
 1536 Böker et al. 2022; P. Jakobsen et al. 2022) and Gemini
 1537 Observatory’s AO assisted GIRMOS (U. Conod et al.
 1538 2023) can address the issue of age-metallicity degener-
 1539 acy in galaxy profiles through spatially resolved spectral
 1540 fitting. Along with MUSE, cutting-edge technologies of
 1541 Atacama Large Millimeter/submillimeter Array, James
 1542 Webb Space Telescope’s MIRI, etc. can help us map the
 1543 dust distribution in galaxies (e.g., L. Lin et al. 2020; B.
 1544 Magnelli et al. 2023; K.-i. Tadaki et al. 2023; J. Li et al.
 1545 2024; C. Lind-Thomsen et al. 2025).

1546 Collectively, these multi-wavelength and spatially re-
 1547 solved perspectives will allow us to move beyond photo-
 1548 metric decomposition and directly test the evolution-
 1549 ary pathways proposed in this work. For the quies-
 1550 cent population in particular, these data will be instru-
 1551 mental in distinguishing between ‘newcomer’ galaxies —
 1552 whose structures are still influenced by their recent star-
 1553 formation history — and aging QGs shaped by interac-
 1554 tions that follow quenching. By tracing age and metal-
 1555 licity gradients into the stellar halos and reconstructing
 1556 galaxy kinematic properties to large radii, we can quan-
 1557 tify the role of ex-situ growth in driving the wavelength-
 1558 dependent structural evolution of both galaxy bulges
 1559 and disks.

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 1573 vey (CLAUDS), which is a collaboration between as-
 1574 tronomers from Canada, France, and China described in
 1575 M. Sawicki et al. (2019). CLAUDS data products can

1576 be accessed from <https://www.clauds.net>. CLAUDS is
 1577 based on observations obtained with MegaPrime/ Mega-
 1578 Cam, a joint project of CFHT and CEA/DAPNIA,
 1579 at the CFHT which is operated by the National Re-
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 1602 nomical communities of Japan and Taiwan, and Prince-
 1603 ton University, USA. The Hyper Suprime-Cam (HSC)
 1604 collaboration includes the astronomical communities of
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1613 AUTHOR CONTRIBUTIONS

1614 This work is a part of the Ph.D. research of AG at
 1615 Saint Mary’s University, Halifax. AG was responsible
 1616 for writing and submitting the manuscript. His Ph.D.
 1617 supervisors, ID and MS were actively involved in the
 1618 data analysis and interpretation. They also obtained the
 1619 funding and edited the manuscript. GD produced the
 1620 photometric catalog used in this work, while LC mea-
 1621 sured the stellar masses of galaxies. The other authors
 1622 contributed to the dataset used in this work.

1623 *Facilities:* Subaru Telescope (HSC), CFHT (Mega-
 1624 Cam)

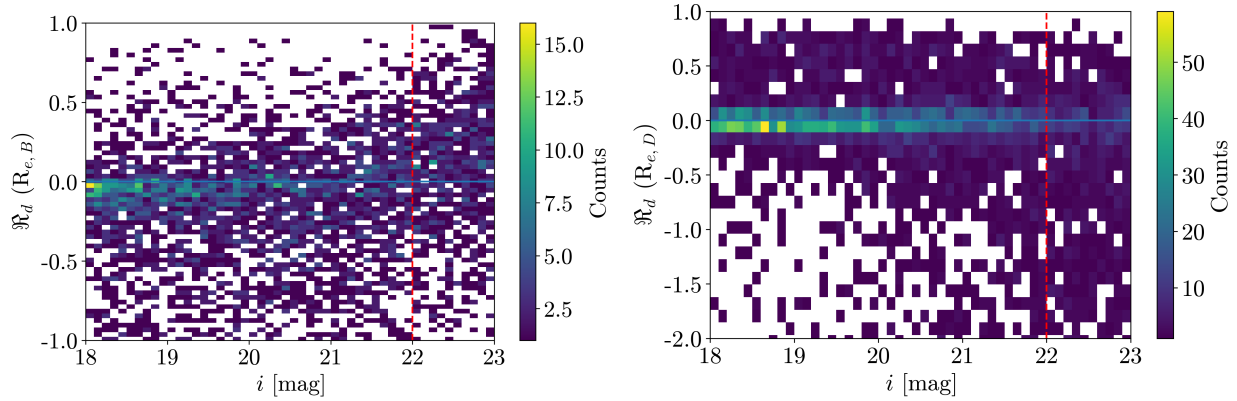


Figure 10. Distribution of relative difference (\mathfrak{R}_d ; Eq. A1) in size measurements of bulges (left) and disks (right) as a function of i -band magnitudes.

1625 *Software:* NUMPY (S. van der Walt et al. 2011; 1630 Hunter 2007), GALFITPYWRAP¹⁶, SExtractor (E.
 1626 C. R. Harris et al. 2020), SciPy (P. Virtanen et al. 1631 Bertin & S. Arnouts 1996), SOURCEEXTRACTOR++ (E.
 1627 2020), ASTROPY (Astropy Collaboration et al. 2013, 1632 Bertin et al. 2020), PSFEX (E. Bertin 2011), DYNesty
 1628 2018, 2022), PHOTUTILS (L. Bradley et al. 2022), SCIKIT- 1633 (J. S. Speagle 2020), LINMIX¹⁷, GALFIT (C. Y. Peng
 1629 LEARN (F. Pedregosa et al. 2011), MATPLOTLIB (J. D. 1634 et al. 2002, 2010)

APPENDIX

A. SIMULATIONS

1636
 1637 To assess the robustness of our SE++ measurements and quantify uncertainties, we conduct simulations by modeling
 1638 and planting mock galaxies into the real CLAUDS+HSC images. Using GALFIT (C. Y. Peng et al. 2002, 2010), we
 1639 generate images of artificial galaxies with a Sérsic bulge and an exponential disk. In doing so, we explore the Sérsic
 1640 parameter space of local galaxies modeled in the literature (L. Simard et al. 2011). We generate 5,000 synthetic galaxies
 1641 with properties uniformly spanning the observed parameter space: bulge effective radii $0.1 \leq r_{e,B}/\text{arcsec} \leq 2.0$, disk
 1642 effective radii $0.5 \leq r_{e,D}/\text{arcsec} \leq 10.0$, Sérsic indices $0.5 \leq n_B \leq 8.0$, and B/T ratios $0.1 \leq B/T \leq 0.9$.

1643 To avoid artificial overcrowding, we only plant 50 simulated galaxies at a time in a CLAUDS+HSC image (\sim
 1644 0.168×0.168 arcmin²). This approach tests the recovery of the bulge and disk parameters under realistic conditions,
 1645 taking into account the effects of noise, crowding, and PSF.

1646 To analyze the simulation results, we explore the relative difference between input and output values of a given
 1647 parameter x , defined as

$$1648 \quad \mathfrak{R}_d(x) = \frac{x(\text{input}) - x(\text{output})}{x(\text{output})}. \quad (\text{A1})$$

1649 In Figure 10, we show the distribution of relative differences in effective radii of bulges (left) and disks (right) as
 1650 a function of galaxy magnitudes in i -band. Although SE++ is able to recover the sizes, the scatter in the relative
 1651 difference increases at $i > 22$ mag. Our simulations also show that we can reasonably model other Sérsic parameters
 1652 for bulge and disk components for galaxies with brightness $i < 22$ mag. Beyond this brightness limit, the noise
 1653 dominates and the uncertainties in the measured parameters are too large. Hence, we limit our analysis only to bright
 1654 ($i < 22$ mag) galaxies in this work.

A.1. Uncertainty Estimation

1655
 1656 To estimate the uncertainties in our fitting results, we follow a method similar to the one prescribed in A. George
 1657 et al. (2024) and compare the input parameters of the simulated galaxies with the fitting results from SE++. To

¹⁶ <https://github.com/Grillard/GalfitPyWrap>

¹⁷ <https://linmix.readthedocs.io/en/latest/src/linmix.html>

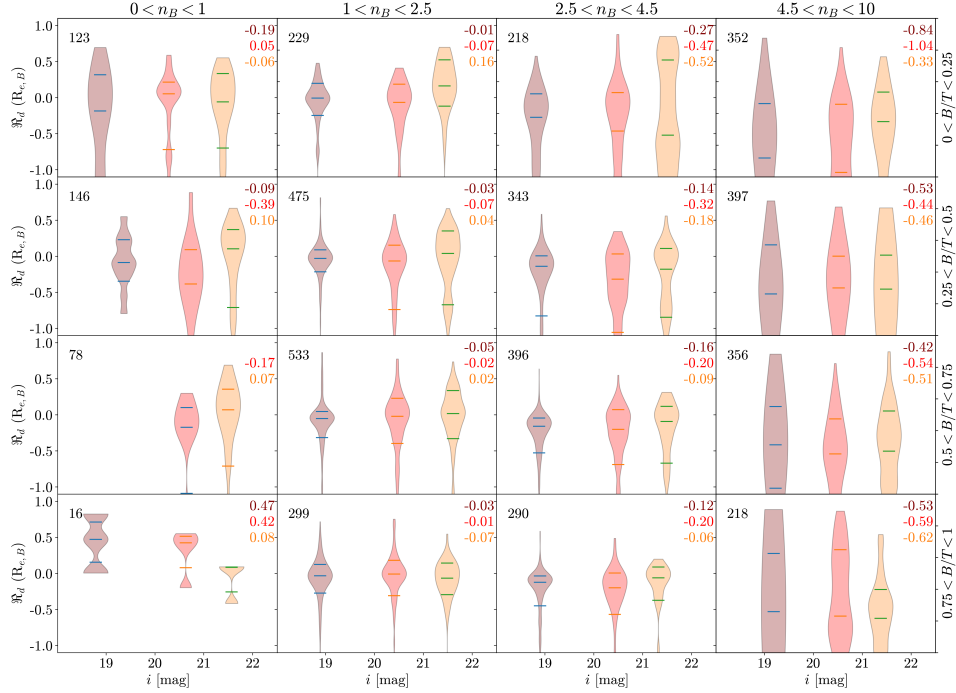


Figure 11. Distribution of relative difference in size measurements of bulges ($\mathfrak{R}_d(R_{e,B})$; Eq. A1) in bins of i -band magnitudes, bulge Sérsic index (n_B) and bulge-to-total ratio (B/T). In each subpanel, the median $\mathfrak{R}_d(R_{e,B})$ of the violin plots is provided in the top-right corner in its respective color. We also provide the total number of simulated galaxies in each subpanel at the top left corner.

1658 quantify the comparison, we use the relative difference in Eq. A1 and investigate how this $\mathfrak{R}_d(x)$ varies with input
 1659 magnitudes and output structural parameters, B/T and n_B for bulge components (Figure 11) or B/T and $R_{e,D}$ for
 1660 disk components (Figure 12). We then adopt the scatter in $\mathfrak{R}_d(x)$ in the three-dimensional bins of magnitude, B/T
 1661 and $R_{e,D}$ or n_B as the uncertainty in the measurement of the fit parameter x .

1662 A.2. Bias Corrections to the Size Measurements

1663 Because our ground based measurements are affected by atmospheric seeing, the simulations also show that there
 1664 can be some systematic biases in the size measurements. For example, we tend to overestimate the sizes of bulges
 1665 if the modeled $n > 4.5$ (right panels in Figure 11). To reduce the impact of these biases in our results, we correct
 1666 for the offsets present in our simulations such that the median values of $\mathfrak{R}_d(R_{e,B})$ and $\mathfrak{R}_d(R_{e,D})$ in the 3-D bins of
 1667 Figures 11 and 12, respectively, are 0. We estimate these bias-corrected effective sizes as

$$1668 R_{e,x}^{\text{corr}} = R_{e,x} (1 + \text{med} [\mathfrak{R}_d(R_{e,x})]_{3\text{D bin}}), \quad (\text{A2})$$

1669 where x is either B (bulge) or D (disk), and $\text{med} [\mathfrak{R}_d(R_{e,x})]_{3\text{D bin}}$ is the median offset in the 3-D bins as shown in the
 1670 simulation violin plots (Figures 11 and 12).

1671 Consequently, the size measurements presented in Section 4 and utilized throughout this analysis represent these
 1672 bias-corrected values. However, we note that the major qualitative results of this work remain valid even if we use
 1673 uncorrected size measurements.

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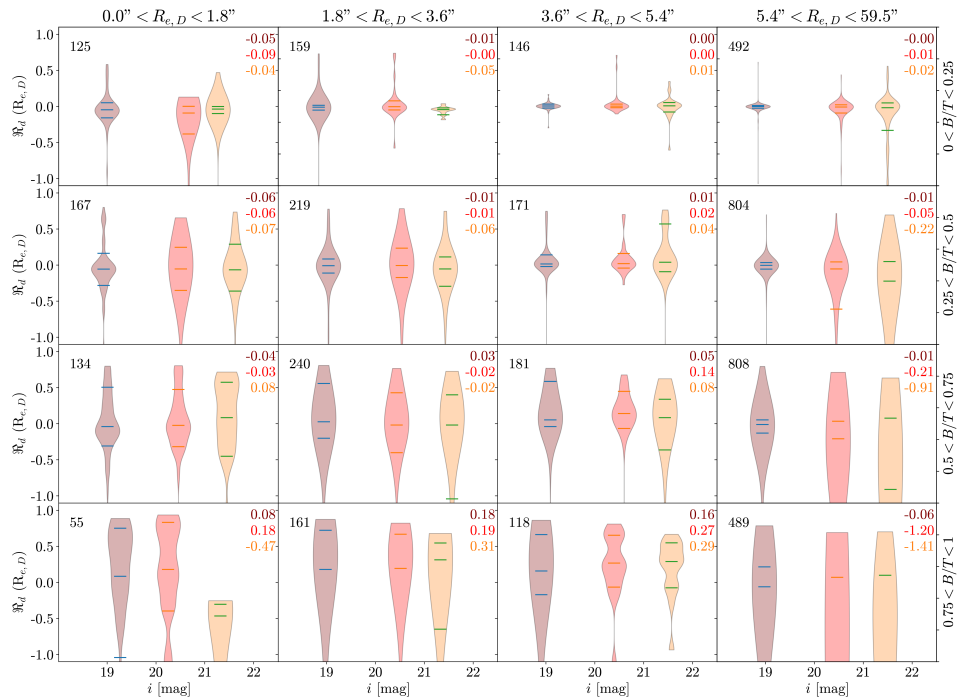


Figure 12. Distribution of relative difference in size measurements of disks ($\mathfrak{R}_d(R_{e,D})$; Eq. A1) in bins of i -band magnitudes, $R_{e,D}$ and B/T . As in Figure 12, the median $\mathfrak{R}_d(R_{e,D})$ values and the total number of simulated galaxies are provided in the top-right and top-left corners of each subpanel, respectively.

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